The photometric effect of errors in Y-position could not be determined directly to better than 5 percent from the observed Pt-Cr-Ne emission line spectra. However, the light loss caused by an incorrect Y-center can be modeled using the mean point source cross-sections and the mean loci of Y-centers. Examples of observed light loss curves derived from high signal-to-noise external Tungsten lamp continuum spectra are presented. These results will provide a useful check on the subsequent modeling.

# 6.3.2 Location of Spectra

# 6.3.2.1 Definition of Y-Base and Theta-Z

In order to determine the center of a spectrum on the photocathode, the diode array is stepped perpendicular to the dispersion to make a Y-map. The standard procedure in ground calibrations was to use 24 Y-steps covering a range of 384 microns, resulting in an image with a step size of 16 microns in Y, and 612 pixels on 50 micron centers in X. The Y-map is centered on the expected position of the spectrum. We are interested in computing two parameters from the Y-map: (1) the optimum Y-deflection for acquiring the spectrum, referred to as the Y-base, and (2) the angle that the spectrum makes relative to the diode array, designated theta-Z.

The optimum Y-deflection is determined in the following manner. Once the Y-center of the spectrum is known as a function of diode number, a linear curve is fit to the points. The data points that define the Y-center can deviate from a straight line because of small distortions in the magnetic focusing fields. The linear fit is then used to define the Y-base. When the

entire diode array is covered by a spectrum, the Y-base is evaluated at diode 256, the X-center of the array. For other dispersers, where only part of the diode array is used, the Y-base is evaluated at the center of the spectrum. These central X-positions are tabulated for all disperser/tube combinations in Table 6.3.2.1-1. Examples of Y-center curves and best fit lines are shown in Figures 6.3.2.1-1 through 6.3.2.1-4. Y-bases are measured in "Y-base units," which correspond to microns to within 10%. The angle theta-Z is defined as the arctangent of the slope of the best fit line.

# 6.3.2.2 Algorithms for Determining the Y-Center of Spectra

Considerable thought was given to the problem of devising a general algorithm to determine the Y-center as a function of X. Since the prime method of determining Y-bases uses the on-board Pt-Or-Ne lamps and since many regions of interest contain few bright emission lines, the algorithm must work well on low signal-to-noise data. Furthermore, many spectra are not well centered in the Y-map, so that the algorithm should be insensitive to truncation of a Y cross-section. The shapes of the oross-sections in Y vary with aperture size from the square shape shown in Figure 6.3.2.2-1 to the triangular form shown in Figure 6.3.2.2-2. Three methods for determining the center of such curves were considered: (1) cross-correlation, (2) centroiding, and (3) contour averaging. The cross-correlation method was selected, because it produced the best results for low signal levels. The merits of each method are briefly discussed in the following sections.

Table 6.3.2.1-1

END POINTS, X-CENTERS, AND EXPOSURE TIMES FOR FOS SPECTRA

GRATING		OINTS ODE)	RED DETECTOR CENTER (DIODE)	EXPOSURE TIME PER Y-STEP (8)
<b>H</b> 13	* `		(52022)	
<b>R</b> 19	1	396	198	. <del>T</del>
H27	ī	512	256	5
H40	ī	512		Ď
<b>H</b> 57	î		256	1
E78		512	<b>2</b> 56	1
	54	512	283	. 1
L15	1	126	63	1
L65	1	214	107	<u>ī</u> .
PRI	833	497	415	0.5
GRATING	ENDP(	OINTS ODE)	HLUE DETECTOR	EXPOSURE TIME
B13	60		CENTER (DIODE)	PER Y-STEP (S)
<b>K</b> 19	•••	512	286	10
H27	<u> </u>	512	<b>2</b> 56	5
	ï	512	<b>2</b> 56	<b>5</b> ·
<b>B40</b>	1	512	256	2
<u> 8</u> 57	1	206	103	5 .
<b>H</b> 78	•	•	•	•
L15	318	512	415	10
<b>L6</b> 5	295	873	834	10
PRI	27	178		10
		1/0	102	0.5

Based on the June 1985 wavelength calibration by Sirk and Bohlin.

<sup>\*</sup>This combination is never used

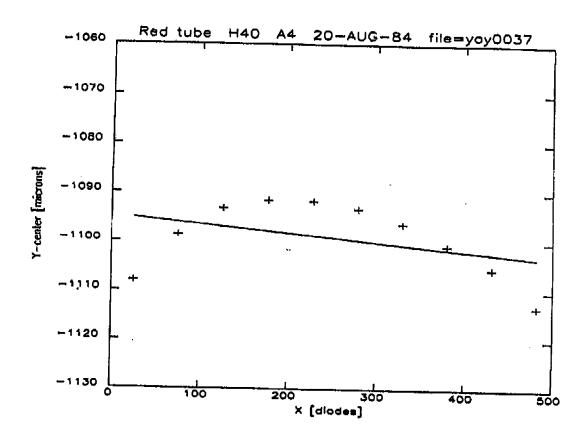


Figure 6.3.2.1-1. Y-center as a function of diode for red tube H4O, ambient calibration, 21 kv voltage, August 20, 1984. Crosses: the Y-center for each bin. The solid line is the least squares fit.

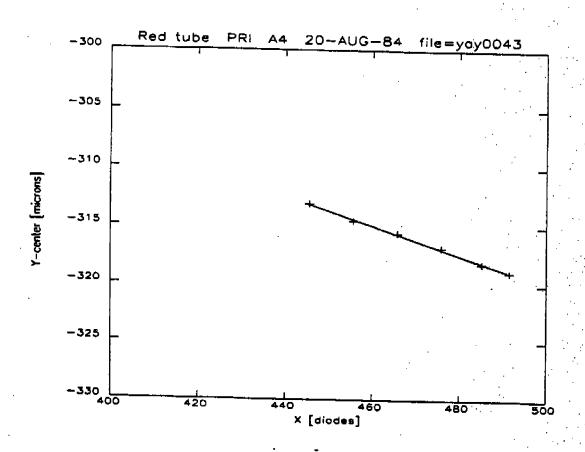


Figure 6.3.2.1-2. Y-center as a function of diode for red tube prism, ambient calibration, 21 kv voltage, August 20, 1984. Note that the X scale is expanded to cover the region from diode 400 to 500 where the prism spectrum is located.

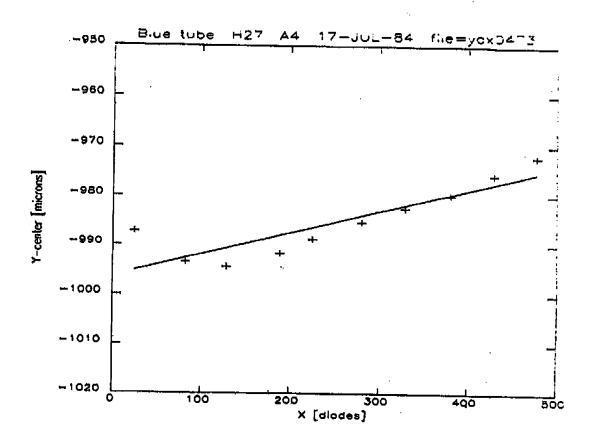


Figure 6.3.2.1-3. Y-center as a function of diode for blue tube H27 vacuum calibration, 23 kv voltage, July 17, 1984.

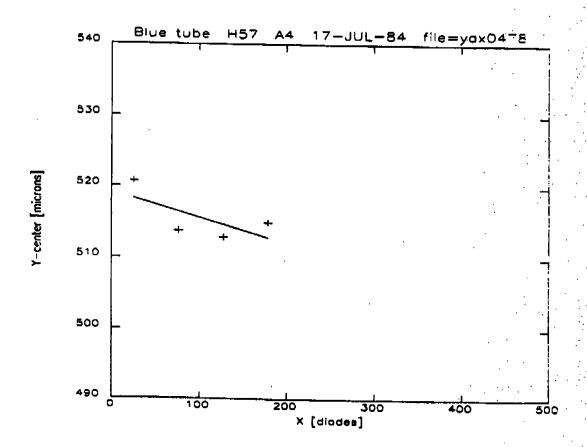


Figure 6.3.2.1~4. Y-center as a function of diode for blue tube H57 vacuum calibration, 23 kv voltage, July 17, 1984. The diode array is only partially covered by the spectrum, because the blue tube attenuates strongly above 5500A.

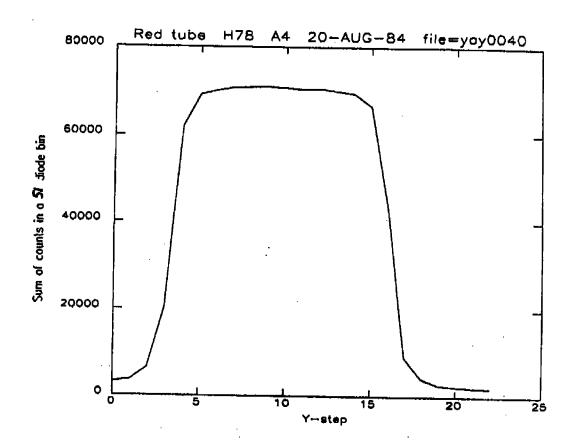


Figure 6.3.2.2-1. Cross-section of the 0.1 arcsec spectrum with a bin size of 50 diodes. The size of the Y-steps is 16 microns.

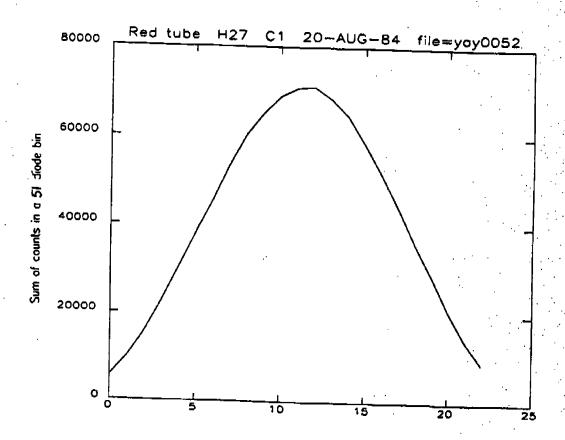


Figure 6.3.2.2-2. Cross-section of an 1.0 arcsec spectrum with a bin size of 50 diodes. The profile is much more rounded than the 0.1 arcsec aperture, because the spectrum is almost as wide as the diode array.

## 6.3.2.2.1 Cross-Correlation

A square template is cross-correlated with the Y crosssection to determine the point of best correlation which defines
the center of the spectrum. In order to improve the counting
statistics, the spectra are summed in 51 diode bins, resulting in
10 data points per spectrum in X. The bin size is a compromise
between the need to produce smooth cross-sections in regions of
the spectrum with few emission lines, and the need for enough
data points in X to measure the curvature adequately. The effective center of each bin depends on the distribution of spectral
lines within it. The bin center is calculated with a centroid
in X:

$$\overline{X} = \frac{\sum X_i C_i}{\sum C_i}$$

where  $x_1$  are the diode numbers and  $c_1$  the corresponding count rates.

To simplify template fitting, each Y cross-section is normalized to its peak. Only two template widths are necessary to accomodate all the aperture sizes, because for apertures one arcsec and smaller the FWHM is determined by the 1.43 arcsec (200 mioron) diede height. The remaining three apertures are 2 arcsec wide in the Y-direction, producing 280 micron cross-sections. A quadratic fit is made to the three points at the peak of the correlation curve to determine the center of the spectrum. The algorithm is the same as that used to determine the line centers in the FOS wavelength calibration.

This method produces satisfactory results when the number of counts in the peak of the cross-section is greater than 200. Results for bins with fewer than this number of counts are not included in the Y-center plots. This is particularly important for the low resolution and prism dispersers where the spectrum is confined to a small portion of the diode array.

## 6.3.2.2.2 Centroiding

The centroiding method is simpler and computationally faster than cross-correlation, but has two major drawbacks: (1) the Y-center is easily thrown off by noise due to poor counting statistics in the cross-section; and (2) centroiding produces incorrect results when a Y-map is truncated, which is often the case in the existing calibration data. A cross-section consists of a set of Y-positions  $y_1$  and the corresponding count rates  $c_1$ . The centroid in Y is then:

$$\overline{Y} = \frac{\sum y_i C_i}{\sum C_i}$$

The truncation problem can usually be alleviated by centroiding only those data points which are above the half-maximum of the cross-sections. The Y-maps are rarely truncated more than this. There is no correction for the errors produced by poor counting statistics in Y-cross-sections, so this method is not considered suitable for general use.

## 6.3.2.2.3 Contour Averaging

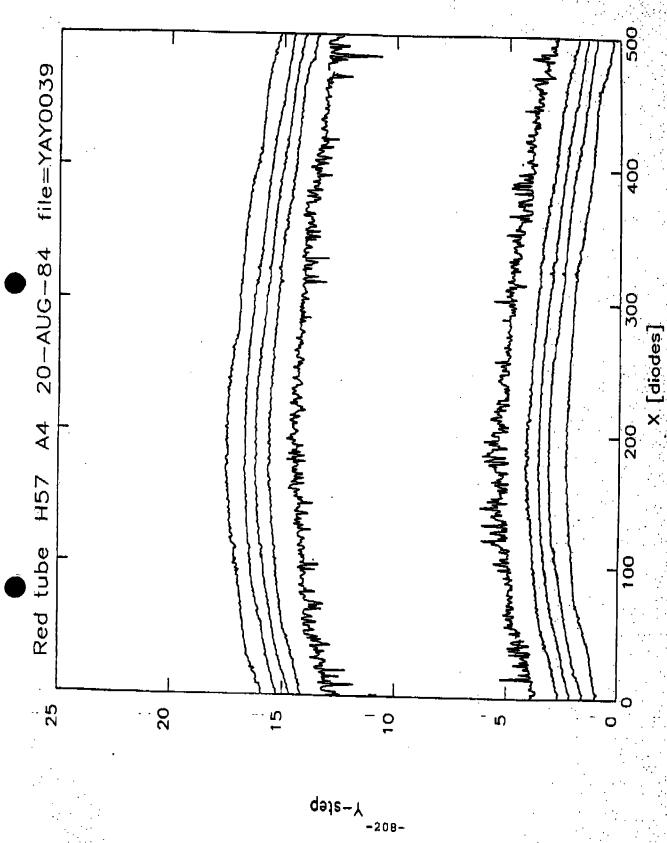
Contour averaging in effect computes the Y-positions of the upper and lower edges of the spectrum, then averages them to find the center. In this method, the contour lines at intensities of

.2, .4, .6, and .8 are computed for a Y-map where each Y crosssection has been normalized to its peak. The contours for a
continuum spectrum are shown in Figure 6.3.2.2.3-1. The average
of the Y-positions of matching contour lines gives the center of
the spectrum in Y along the dicde array. Increasing the number
of contours in the average does not result in more accurate Ycenters, because the additional contours are interpolated between
the same data points in the cross-section.

Contour averaging uses only the data in two or three Y-steps at the upper and lower edges of the spectrum. The cross-correlation method is also influenced mainly by the points at the edges of the spectrum, because that is where the difference between the template and the spectrum cross-section is greatest. Thus, the two methods are quite similar for high signal-to-noise data. Indeed, the two methods agree to within 2 microns (one-eighth the Y-step size), which is remarkably good agreement considering the 200 micron resolution of the Y-maps. The main drawback to contour averaging is that it does not work well on noisy cross-sections, because linear interpolation is used to calculate the contours, and interpolation does not produce satisfactory results on non-monotonic functions. Therefore, cross-correlation is preferred over the contouring technique, since cross-correlation gives more precise results with low count-rate data.

### 6.3.2.3 Y-Base and Theta-Z Measurements

Numerous Y-base and theta-Z measurements were made during the vacuum calibration of July 1984 and the ambient calibrations of June and August 1984. The vacuum Y-maps were made with the



trum, clearly showing the curvature produced by the red Digicon. The cona continum spec-117 -map of Y Y A contour plot of a normalized ! 155 ы Figure 6.3.2.2.3-1.

Pt-Cr-Ne emission line lamp, while the ambient Y-maps use the Tungsten and Deuterium external continuum lamps. The spectra were corrected for paired pulse effects with the constants  $t_1 =$ -0.17  $\times$  10<sup>-6</sup> and  $t_2 = 10.5 \times 10^{-6}$  seconds. All Y-base measurements taken since the removal and replacement of the Digicons in the spring of 1984 are summarized in Table 6.3.2.3-1. The range referred to in this table is the total variation in Y-center across the diode array, which is affected both by theta-Z and the amount of distortion. The FOS file-name listed in sixth column consists of a three letter tape designation followed by the number of the file. Files on a given tape are numbered consecutively, e.g., file YAA0001 is the first file on tape YAA. Not all consecutive files were recorded without filter-grating wheel movement. The scatter in the Y-base values due to the filtergrating wheel repeatability is discussed by Hartig, Bohlin, and Harms in CAL/FOS-012 and by Hartig in CAL/FOS-017.

The value of theta-Z is highly dependent on where the endpoints of the spectrum are chosen for the low dispersion modes
with short spectra, and for high dispersion modes with partial
coverage (H19, H78 Red and H13, H57 Blue). For example, on H57
blue, the theta-Z derived from the first 200 diodes is -0.035
compared with +0.035 if fainter spectral lines out to diode 350
are used. However, the calculated Y-bases differed by only three
microns. Because the faint lines are not detected in short
exposures, the endpoint of the spectrum was set at diode 200.
The typical error in the determination of theta-Z is .01 degrees
for gratings that utilize the entire diode array, and 0.5 degrees

where theta-Z is calculated from fewer data points. The Y-base values are accurate to better than 8 microns in all cases. Since the Y-base depends on where the endpoints are chosen, the endpoints of the spectra used to compute theta-Z are listed in Table 6.3.2.1-1. Exposure time may also affect the measured value of theta-Z. Underexposure may result in the loss of fainter spectral lines at the ends of a spectrum. To avoid this, we recommend using the optimum exposure times for the .01 arcsec aperture listed in Table 6.3.2.1-1.

#### 6.3.2.4 Shifts Due to Changes in Temperature and Digicon Voltage

The Y-base and theta-Z measurements were examined to investigate whether systematic shifts occurred between data taken at different temperatures. There are three temperatures to be considered: cold operate in vacuum at -30°C, hot operate in vacuum. at -10°C, and ambient calibration at 20°C. All observations at a given temperature were averaged in this analysis. On the red tube, there was no systematic shift in Y-base due to temperature: the mean Y-base shift between cold operate and ambient was -17 microns with a standard deviation of 34 microns. However, the red tube did show a systematic change in theta-Z of .052  $\pm$  .012 degrees between cold operate and ambient. The blue tube showed a systematic shift in Y-base of -61 + 31 microns between cold operate and ambient, but the blue side theta-Z did not show a systematic shift: the mean was shift -.004  $\pm$  .026 degrees. The offsets due to changes in temperature are listed in Tables 6.3.2.4-1 and 6.3.2.4-2.

Table 6.3.2.3-1
Y-BASE AND THETA-Z MEASUREMENTS

<b>GRATING</b>		Y-BASE (microns)	BANGE IN Y (microns)	THETA-Z (degrees)	PILE	MAX OTS
RED TUBE,	AMBIENT,	22 JUNE 1	. <b>984, TEMP=R</b> 0	OM, V=18KV,	Tungsten	LAWP
H27	.44	<b>3</b> 15	22	-0.032	YAQ0097	1372
<b>E27</b>	A4U	751	24	-0.045	YAQOO9B	1270
H40	<b>A4</b>	-1122	25	-0.023	YAQ0078	1076
H40	A4	-1122	25	-0.022	YAQ0077	1090
<b>H40</b>	a4u	~687	24	-0.020	YA00078	680
<b>H40</b>	A4U	-681	21	-0.014	YA00079	675
<b>H4</b> 0	A4U	-698	28	-0.021	7A00080	685
<b>H</b> 57	<b>A4</b>	-1313	32	-0.052	YAQ0082	2567
<b>3</b> 57	<b>A4</b>	-1313	32	-0.051	YAQ0088	551
<b>H57</b>	A4 ·	-1281	. 27	-0.043	YAQ0090	1320
H57	A4	-1281	31	-0.049	YAQ0091	1137
E57	<b>A4</b>	-1296	32	-0.051	YAQ0092	1074
H57	<b>A4</b>	-1297	32	-0.051	EROODAY	1072
<b>H57</b>	A4U	-863	31	-0.049	YAQ0094	1011
H57	<b>B</b> 2	-1105	31	-0.050	YAQ0107	5876
H78	A4U	628	21	0.006	YAQ0087	1376
H78	<b>A4</b>	186	21	0.006	<b>YAQOO88</b>	2262
1178	<b>A4</b>	186	20	0.004	YAQ0089	1125
L85	<b>A4</b>	-361	87	0.23	YAQOO99	23350
<b>1</b> 65	<b>A4</b>	-360	- 87	0.24	YAQ0100	4790
L65	<b>A4</b>	-359	87	0.23	YAQ0101	4680
<b>L</b> 65	A4U	77	87	0.24	YAQ0102	2930
PRI	14	-872	5	-0.12	YAQ0104	55158
PRI	A4U	71	7	-0.15	YAQ0105	33629
BLUE TUBE	AVBIRNT	, 22 June	1984, TEMP=R	00M, V=18KV	, Tukuste:	N LAMP
<b>H40</b>	A2	544	81	0.086	YAR0022	1840
<b>H</b> 57	<b>A4</b>	<b>58</b> 5	10	0.039	YAROO28	597
L65	<b>A4</b>	-203	26	0.21	YAR0024	1172
elue tube	, VACUUM,	11 JULY 1	984, TEMP=-3	00, <b>V=23K</b> Y,	PLATINUM	LAMP
<b>H</b> 18	<b>A4</b>	-249	13	-0.037	YAT1896	1427
H13	<b>A4</b>	-249	15	-0.039	YAT1897	1437
<b>H</b> 13	44	-261	17	-0.044	YAT1898	1953

GRATING	APERTURI	Y-RASE (microns)	EANGE IN Y (microns)		FILE	MAX OTS
		-	•,	(degrees)		
BLUE TUBE	, VACUUM,	11 JULY	1984, <b>TEM</b> P≃-3	00, V=23KV	, PLATINUM	LAMP
H19	. A4	-536	19			- '
H19	Å4	-538	22	0.088	YAT1876	87
E19	Ä4	-537	24	0.10	YAT1877	89
		-001	24	0.11	YAT1878	<b>7</b> 5
E27	<b>A4</b>	-981	25			
B27	Ã4	-972	25	0.064	YAT1874	553
			20	0.060	YAT1875	527
<b>E40</b>	<b>A4</b>	520	16	0.040	<b></b>	
<b>B40</b>	Ā4	514	19	0.043	YAT1872	1732
			19	0.055	YAT1873	1994
E57	A4	<b>5</b> 56	17			
			11	-0.044	YAT1899	1953
L65	A4	-314	24			
L65	A4	-319	34	0.14	YAT1884	1208
L85	Å4 .	-319	29	0.20	YAT1885	1235
L65	A4	<b>-267</b>	25 25	0.17	YAT1885	1160
L65	Ã4	-283	29	0.15	YAT1890	1203
		-200	AV	0.18	YAT1894	1182
PRI	<b>A4</b>	-390	6			
PRI	Ā4	-393	5	-0.11	YAT1887	1774
PRI	Ā4	-341		-0.09	YAT1888	2023
PRI	. A4	- <b>3</b> 51	2	-0.01	YAT1892	2061
			3	-0.06	YAT1895 .	2166
RED TUBE.	VACUUM.	15 JIT.V 10	94 <b>8773</b> 70 004		_	
•	,	4021 19	84, TEMP=-300	), Y=21KY, }	PLATINUM I	AMP
H19	<b>A4</b>	-218	83			
H19	Ā4U	216	35 35	-0.077	YAX0020	2854
			90	-0.083	YAX0021	2405
H27	<b>A4</b>	316	48			
<b>H40</b>	<b>A4</b>	-1112	36	-0.11	YAX0022	10950
<b>H</b> 57	· A4	-1207	45	-0.082	YAX0025	6948
H78	A4U	567	18	-0.10	AYX0058	<b>3133</b> 5.
L15	Ā4	-121	0	-0.058	YAX0029	14333
			-	-0.10	XXX0030	400
BLUE TURE	VACUUM.	17 JULY 16	984, <b>TEM</b> P=-10	VI. Do Lar		
			and Inter-10	~, <b>γ=23</b> Εγ,	PLATINUM )	
H13	A4U	+214	19	-0.040	-	
		<b>-</b>		-0.040	YAX0487	7800
H19	<b>A4</b>	-514	28	A 070	***	-
H19	<b>A4</b>	-499	28	0.070	YAX0474	8500
H19	A2	-501	25	0.084	YAX0488	80
H19	<b>C</b> 1	-501	25	0.067	YAXO486	2500
				0.068	YAXO489	7000

GRATING	APERTURE	Y-BASE (microns)	RANGE IN Y (microns)	THETA-Z (degrees)	PILE	MAX OTS
BLUE TUBE	, VACUUM,	17 JULY	1984, TEMP=-1	100, <b>V=23kY</b> ,	PLATINUM	LAMP
H19	<b>B</b> 2	-289	25	0.067	YAX0481	2097
H19	<b>B</b> 3	-287	25	0.088	YAXO482	6209
H19	A2	-501	26	0.085	YAXO486	2602
H19	A3	-502	82	0.074	YAX0487	1995
H19	<b>01</b>	-501	24	0.083	YAX0489	7027
<b>H</b> 19	02	-300	27	0.065	YAX0490	8743
<b>H27</b>	<b>A4</b>	-985	22	0.050	YAX0478	18000
H40	A4	507	16	0,037	YAX0476	25000
R40	<b>A4</b>	500	15	0.034	YATOSOO	4007
<b>H40</b>	A4U	923	16	0.037	YAX0499	4152
R40	01	499	15	0.038	YAXO498	81903
<b>H40</b>	Ç1	499	14	0.037	YAX0497	31994
<b>H40</b>	Q1	920	14	0.036	YAXO498	31684
E40	ea.	497	15	0.040	YAX0501	14087
H40	USA	916	15	0.036	YAX0502	13928
<b>H40</b>	a2U	920	15	0.037	<b>EOGOXAY</b>	10385
<b>H40</b>	<u> A2</u>	508	15	0.039	YAX0504	10438
<b>H40</b>	<b>B</b> 3	714	15	0.038	<b>YAX05</b> 05	28848
<b>H40</b>	<b>B</b> 2	711	16	0.042	YAX0508	<b>D</b> 882
<b>H40</b>	<b>B</b> 1	714	15	0.039	YAX0507	9719
H57	<b>A4</b>	<b>51</b> 5	9	-0.043	YAXO478	20000
L15	A4U	+6	2	0.008	<b>YAX0469</b>	2000
L65	14	-274	21	0.22	YAX0471	24000
RED TURE,	VACUUM,	17 JULY	1984, TEMP=-10	00, Y=21KŸ,	PLATINUM :	LAMP
R19	<b>A4</b>	-212	82	073	YAX0522	2600
<b>E27</b>	A4	+342	42	094	YAXO523	30000

GRATING		Y-BASE (microns)	PANGE IN Y	THETA-Z (degrees)	FILE	MAX CTS
BLUE TUBE	, VACUUM,	20 JULY	1984, <b>TEM</b> P=-1	.oc, <b>v=23k</b> v,	PLATINUM	LAMP
H13 H13	A4	-222	16	-0.094	YAX0902	274
H13	A4U	201	18	-0.11	E000XAY	824
H13	A4	-234	20	-0.12	YAX0920	
ита	A4U	184	13	-0.081	YAX0921	271
RED TUBE	, VACUUM,	20 JULY	1984, <b>TEM</b> P=-1	00, V=21KV,	PLATINUM	LAWP
H78	A4	151	37	-0.12	YAX0907	9898
<b>278</b>	A4	169	25	-0.078	YAXOOOB	9338
E78	A4U	<b>6</b> 06	22	-0.072	YAXOOOS	6698
	AMBIENT,	20-21 AU	GUST 1984, TI	EMP=ROOM, V=	21kV, TÜN	gsten lamp
H19	01	<b>-26</b> 0	83	-0.044	<b>YAY0064</b>	323
H19	<u>C3</u>	-40	27	-0.032	YAYOOB2	643
E19	<b>B</b> 3	-57	26	-0.035	YAY0085	248
H27	<b>A4</b>	289	32	-0.056	YAY0047	351
H27	A4U	714	30	-0.052	YAYOO36	250
H27	A2	290	32	-0.056	YAY0045	
R27	Ã3	290	82	-0.057	YAY0046	6181 1460
H27	B1	489	81	-0.054	YAY0048	4798
H27	<b>B</b> 2	492	80	-0.052	YAYOO49	1673
H27	<b>B</b> 3	491	81	-0.056	YAY0050	9854
<b>E</b> 27	C1	288	33	-0.059	YAY0052	12616 :
H27	C2	508	31	-0.058	YAY0053	9000
<b>E</b> 27	C3	510	29	-0.054	YAY0054	24039
H27	<b>04</b>	502	84	-0.082	YAY0055	20810
740			_			
H40	<b>A4</b>	-1099	20	<b>-0.02</b> 0	YAY0037	4500
<b>H</b> 57	<b>A4</b>	-1265	82	-0.048	<b>YAY0039</b>	10000
<b>H</b> 78	<b>A4</b>	140	23	0.002	YAY0058	11088
178	Ā4U	593	23	0.007	YAY0040	7000
H78	A2	140	22	0.002	YAY0056	54946
E78	A3	142	22	0.002	YAY0057	
<b>H</b> 78	B1	342	21	0.002	YAY0059	
H78	<b>B</b> 2	844	20	0.003	YAY0060	41557 82906
H78	· C2	358	21	0.006	YAY0061	
		7-0	W.L	V.000	TEICOST	40127
L65	<b>A4</b>	-360	28	0.23	YAY0044	5000
PRI	<b>A4</b>	-317	5	-0.043	YAY0043	60000
			_	040	***********	•••••

GRATING	Aperture (		RANGE :	IN Y THET.		le nax cts
BLUE TUBE,	AMBIENT,	26 AUGUST	1984,	TEMP=ROOM,	Y=23XV,	Tungsten la <u>n</u> p
H10 H27 H40 H57 L65 PRI	01 A3 A4 A4 A4 A4	-561 -1057 412 476 -336 -434	81 28 20 7 50 3	0.07 0.08 0.08 -0.01 0.16 -0.04	72 YAZ 64 YAZ 64 YAZ 63 YAZ 7AZ	0106 327 0107 1128 0108 806 0109 1130 0110 2038

The lower half of a paired aperture was used unless otherwise noted with a U for the upper aperture.

#### Table 6.3.2.4-1

## Y-BASE OFFSETS DUE TO CHANGES IN TEMPERATURE

## Red Tube , V=21kv, 0.1 arcsec aperture

Grating E19	Y-base T=-30 (microns) -218	Y-base T=-10 (microns) -212	Shift between T=-30 and t=-10 ) (microns) 6	Y-base Ambient (microns) -261	Shift between T=-30 and ambient (microns) -43
<b>H27</b>	<b>3</b> 16	342	26	289	-27
<b>H4</b> 0	-1112			-1099	ĺŝ
<b>1</b> 57	-1207			-1265	-52
<b>H78</b> U	567			593	26
				1	lann

# Blue tube , V=23kv, 0.1 arcsec aperture

Grating H19 H27 H40 H57 L65 PRI	Y-base T=-30 -537 -977 517 556 -300 -390	Y-base T=-10 -514 -985 507 516 -274	Shift between T=-30 and t=-10 23 -8 -10 -40 44	Y-base Ambient -561 -1057 412 477 -336 -434	Shift between T=-30 and ambient -24 -80 -105 -79 -38	5
	•		Mean: 2 +/-32		Mean: -61 +/- 81	

#### Table 6.3.2.4-2

## THETA-Z OFFSETS DUE TO CHANGES IN TEMPERATURE

## Red Tube , V=21kv, 0.1 arcsec aperture

Grating	Theta-Z T=-30 (degrees)	(degrees)	N=-30 and t=-10 (degrees)	Theta-Z Ambient (degrees)	Shift between T=-30 and an (degrees)	
Ħ19	077	072	.005	044	.033	
H27	109	093	.016	056	.053	
<b>H40</b>	082			020	.082	
<b>H</b> 57	100			048	052	
<b>H</b> 78U	058	043	.015	.007	.085	• , :
					ean: .053 +/-	.012

#### Blue Tube , V=23kv, O.1 arcsec aperture

Grating E19 E27 E40 H57	Theta-Z T=-30 .106 .062 .049 044	Shift between T=-30 and t=-10026012015 .001	7 Ambient 7 .072 .055 .054014	Shift between T=-30 and ambien 084 007 005 .030	t
•		Mean: .013 +/-	.011 Nes	a:004 + /02	6

## Table 6.3.2.4-3

# Y-BASE OFFSETS DUE TO CHANGE IN VOLTAGE

# Red Tube, ambient, 0.1 arcsec aperture

E27 E27U E40 E57 E78 E78	Y-base 18kv (microns) 315 751 -1044 -1226 188 827	Y-base 21kv (microns) 289 714 ~1099 ~1265 140 593	Shift between 18kv and 23kv -26 -37 -55 -39 -48 -34
			Mean: -40 4/- 10

# Blue Tube, ambient, 0.1 arcsec aperture

Grating E19 E27 E40 E57 L65 PRI	Y-base 18ky (microns) -437 -954 544 585 -203 -295	Y-base 23kv (microns) -561 -1057 412 476 -336 -434	Shift between 18kv and 23kv -124 -103 -132 -109 -133 -139
			-109 ean: -123 +/- 14

## Table 6.3.2.4-4

# THETA-Z OFFSETS DUE TO CHANGE IN VOLTAGE

## Red Tube, ambient, 0.1 arcsec aperture

Grating	Theta-Z 18ky	Theta-Z 21kv	Shift between
H27U H40 H57 H78 H78U	(degrees) 045 012 049 006 .006	(degrees) 052 020 048 .002 .007	18kv and 23kv 007 008 .001 004 .001

# Blue Tube, ambient, 0.1 arcsec aperture

Grating	Theta-Z 18kv	Theta-Z 23kv	Shift between
	(degrees)	(degrees)	18kv and 23kv
H40	.086	.054	032
H57		~.013	052

During the July vacuum calibration and the August ambient calibration, the blue and red Digicons were operated at 23 and 21 kV, respectively. However, in the June ambient calibration, data was obtained at 18 kV for both detectors. Both detectors showed systematic decreases in Y-base values with increased voltage. The mean Y-base shift for the red tube between the June and August calibrations was  $-40 \pm 10$  microns. The blue tube Y-base measurements from June are only available on tape for gratings H40, H57, and L65. This data was supplemented by Y-base values calculated in FOS Calibration Notebook 3 for gratings H19, H27, and the prism. The mean shift for these six dispersers is -123  $\pm$  14 microns. When Y-base values from the FOS notebook are compared to those calculated from the data tapes by the method discussed here, the results agree to within 5 microns. Theta-Z did not shift systematically with voltage on the red tube: the mean shift was .007  $\pm$  .009 degrees. Theta-Z comparisons between 18 and 23 kV are available for only two gratings on the blue tube. The shift on H40 was -.032 degrees, and -.052 degrees on H57. Theta-Z comparisons were made using only those gratings with full spectral coverage. The offset due to changes in voltage are shown in Tables. 6.3.2.4-3 and 6.3.2.4-4.

## 6.3.2.5 Mean Point Source Cross-Section

The main interest in FOS photometric precision is for stellar point sources, which are adequately approximated by the fully illuminated 0.1 arcsec aperture. The mean cross-section of an FOS spectrum taken with the 0.1 arcsec aperture was computed by averaging ten cross-sections from one spectrum on each detector.

Each cross-section is the sum of the counts in a 50-diode bin. The cross-sections were normalized to their peaks, and then interpolated on a one micron grid. The center of each crosssection was determined by cross-correlation, and the interpolated cross-sections were then centered and summed to produce the mean curve. The H76 Y-map in file YAY0040 was chosen for the red tube cross-section, because it was one of the few that was not truncated, and yet had sufficient counts for one percent accuracy. The H40 Y-map in file YAX0501 was used for the blue tube cross-section. The FWHM is 206.0 Y-base units (equivalent to microns to within 10%) for the blue curve and 206.3 Y-base units for the red curve. The mean cross-sections are tabulated in Tables 6.3.2.5-1 and 6.3.2.5-2, and plotted in Figures 6.3.2.5-1 and 6.3.2.5-2, where Y-base units and microns are used interchangeably, even though Y-base units are only approximately equivalent to microns.

#### 6.3.2.6 Mean Reduced Locus of Y-Centers

The average Y-center as a function of diode number in an FOS spectrum was determined for each detector. The effect of the different orientation (theta-Z) of each grating was removed by computing the deviation of the Y-center from the best fit line for each grating. The results for the three fully covered gratings on each side (Blue - H19, H27, H40; Red - H27, H40, H57) were shifted to zero Y-base and averaged to obtain the final curves, illustrated in Figures 6.3.2.6-1 and 6.3.2.6-2, and listed in Table 6.3.2.6-1.

Table 6.3.2.5-1
BLUE TUBE MEAN POINT SOURCE CROSS-SECTION

Y			. * *.
(MICRONS)	INTENSITY	Y	
-150.00	4.13142E-02	(MICRONS)	Intensity
-145.00	4.49599E-02	0.0000	1.0000
-140.00	5.01553B-02	5.0000	0.99989
-135.00	5.93092B-02	10,000	0.99887
-130.00	7.07596E-02	15.000	0.99739
-125.00	9.03191B-02	20.000	0.99590
-120.00	0.13048	25.000	0.99490
-115.00	0.18436	35.000	0.99354
-110.00	0.27366	40,000	0.99268
-105.00	0.43243	45.000	0.99170
-100.00	0.60533	50.000	0.99061
-95.000	0.77064	55.000	0.98852
-90.000	0.86509	60.000	0.98574
-85.000	0.92411	65.000	0.98217
-80.000	0.96145	70.000	0.97501
-75.000	0.97578	75.000	0.96533
-70.000	0.98225	80.000	0.95167
<b>-65.000</b>	0.98720	85.000	0.91036
-60.000	0.99058	90.000	0.83798
-55.000	0.99310	95.000	0.74179
-50.000	0.99531	100.00	0.59572
-45.000	0.99691	105.00	0.42942
-40.000	0.99758	110.00	0.27304
-35.000	0.99803	115.00	0.16328
-30,000	0.99833	120.00	0.11605
-25.000	0.99826	125.00	8.39821E-02
-20.000	0.99798	130.00	6.39523E~Q2
-15.000	0.99768	185.00	5.40057E-02
-10.000	0.99816	140.00	4.62675E-02
-5,0000	0.99905	145.00	4.02842E-02
		150,00	3.75334E-02

Table 6.3.2.5-2

BLUE TUBE MEAN POINT SOURCE CROSS-SECTION

Y		Y	
(MICRONS)	intensity		
-150.00	2,18572E-02	(MICRONS)	intensity
-145.00	3.04461E-02	0.00000E+00	1.0000
-140.00	4.27927B-02	5.0000	1.0015
-135,00	5.86833E-02	10.000	1.0026
-130,00	8.58453E-02	15.000	1.0027
-125.00	0.12515	20.000	1.0018
~120.00	0.17346	25.000	1.0008
-115.00	0.25360	30.000	1.0003
-110.00	0.35651	35.000	0.99978
-105.00	0.47012	40.000	0.99926
-100,00	0.58599	45.000	0.99859
-95.000	0.69565	50.000	0.99881
-90.000	0.79401	55.000	0.99912
-85,000	0.87521	60.000	0.99734
-80.000	0.92664	65.000	0.99160
-75.000	0.95811	70.000	0.98287
-70,000	0.97718	75.000	0.96440
-65,000	0.98364	80.000	0.92870
-60,000	0.98824	85.000	0.87438
-55,000	0.99301	90,000	0.80381
-50.000	0.99483	95,000	0.70301
-45.000	0.99521	100,00	0.58884
-40.000	0.99533	105.00	0.48980
-35.000	0.99587	110.00	0.36138
-30.000	0.99640	115.00	0.26733
-25.000	0.99654	120.00	0.18314
-20.000		125.00	0.12813
-15.000	0.99867	130.00	9.20825B-02
	0.99705	135.00	6.49018E-02
-10.000 -5.0000	0.99779	140.00	4.71736E-02
-5.000	0.99879	145.00	3.67673E-02

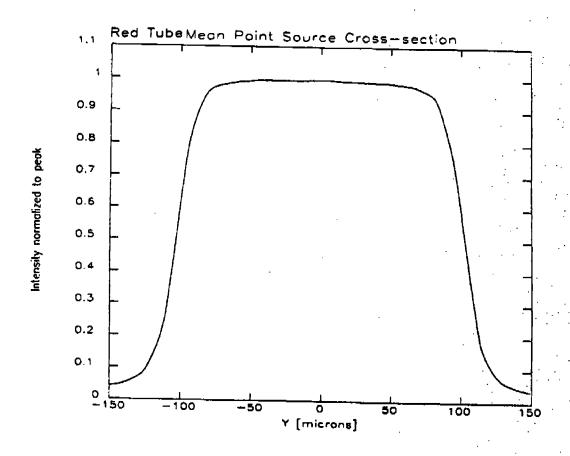


Figure 6.3.2.5-1. Mean cross-sections of a point source spectrum as represented by a 0.1 arcsec aperture spectrum with the red and blue detectors. This cross-section should adequately represent the actual cross-section for a point source in any FOS aperture.

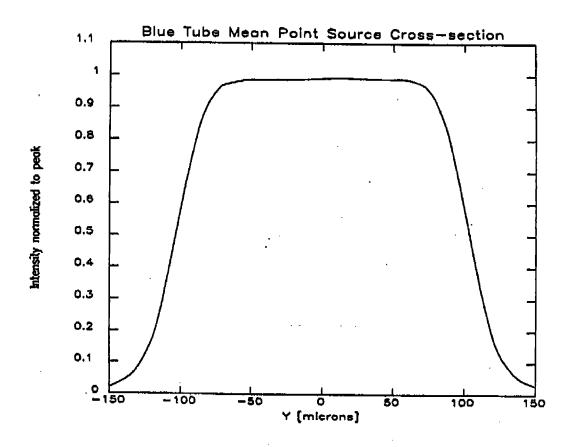


Figure 6.3.2.5-2. Mean cross-sections of a point source spectrum as represented by a 0.1 arcsec aperture spectrum with the red and blue detectors. This cross-section should adequately represent the actual cross-section for a point source in any FOS aperture.

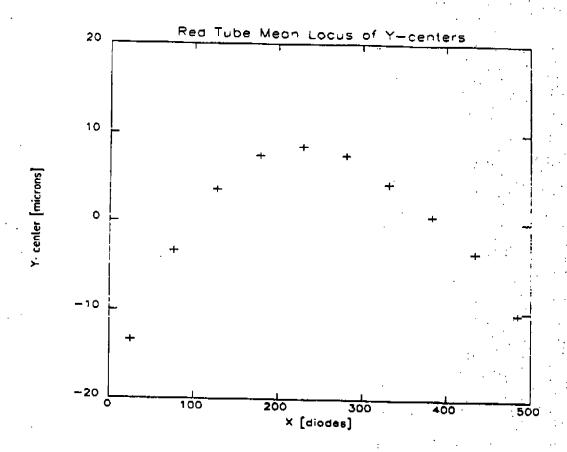


Figure 6.3.2.6-1. Mean Y-center as a function of diode for each detector. The curve has been reduced to zero theta-Z and zero Y-base.

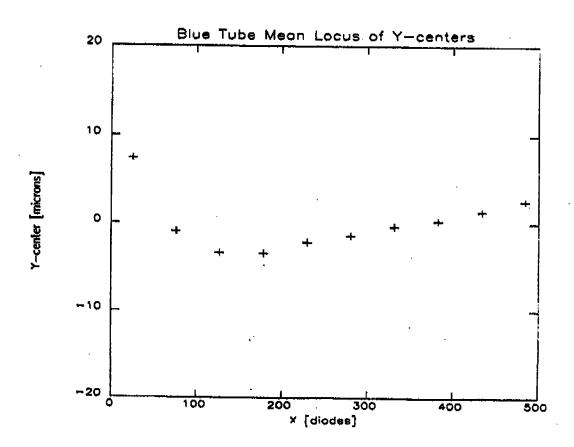


Figure 6.3.2.6-2. Mean Y-center as a function of diode for each detector. The curve has been reduced to zero theta-Z and zero Y-base.

# Table 6.3.2.6-1 MEAN LOCUS OF Y-CENTERS

BLUE TUB	8
DIODE	Y-CENTER
	(MICRONS)
<b>2</b> 5.	7.5
76.	-0.9
127.	-3.4
178.	-3.4
229.	<b>-2.2</b>
280.	-1.4
331.	-0.3
<b>382</b> .	0.3
433.	1.3
<b>484</b> .	2.5

Y-center
(MICRONS)
`-13.3
-3.3
3.6
7.5
8.5
7.5
4.3
0.7
-8.4
-10.2

# 6.3.3 Photometric Consequences of Positioning Errors (Ob-

The photometric consequences of positioning errors must be investigated separately for each combination of detector, grating, and aperture, because theta-Z is different for each grating. Because the locus of Y-centers is not symmetric in Y, a positive Y-base error has a different effect than a negative Y-base error of the same magnitude.

Ideally, we would like to directly measure the light loss, from ground-based calibration data obtained in vacuum. Unfortunately, the existing internal Pt-Cr-Ne emission line spectra in Y-maps simply do not have enough counts to determine the light loss to better than 3-5%. A typical light loss curve derived from an emission line spectrum is shown in Figure 6.3.3-1, illustrating that the curves are noisy and of little use. Some of the continuum spectra taken in ambient calibrations have adequate counting statistics, but because theta-Z changes with temperature, the ambient light loss curves cannot be used to predict the performance of the FOS in vacuum. However, it is possible to predict light loss curves from the mean locus of Y-centers, the mean spectrum cross-section, and the Y-base and theta-Z values measured in vacuum.

Light-loss curves obtained from ambient continuum spectra can be used to check whether the calculated light loss agrees with the observations. The raw data for the light loss measurements are the same Y-maps used for the Y-bass measurements. The Y-maps are summed in 50 diode bins, to improve counting statistics. The spectrum at the nominal Y-base is linearly interpo-

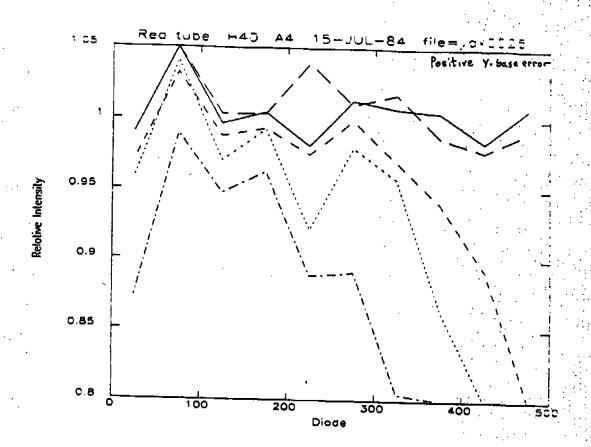
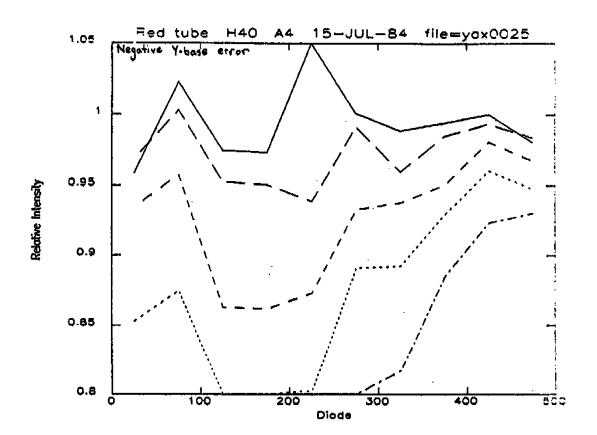


Figure 6.3.3-1. The effects of error in Y-position as a function of diode number determined from an emission line spectrum on the red tube. The upper plot is for positive Y-base errors, the lower plot for negative Y-base errors. The two plots differ because the locus of Y-centers is not symmetric. The five curves represent errors of 30 microns (solid line), 50 microns (long dash), 70 microns (medium dash), 80 microns (short dash), and 90 microns (dot-dash). The error in the curves is at the 5% level.



lated between the Y-steps of the Y-map, where necessary. Similarly, the spectra at offsets of 30, 50, 70, 80, and 90 microns above and below the nominal Y-base are also interpolated. Each light-loss curve at a given offset is simply the binned spectrum at the offset divided by the binned spectrum at the nominal Y-base. Examples of these observed light loss curves are presented in Figures 6.3.3-2 through 6.3.3-4.

### 6.3.4 Photometric Consequences of Positioning Errors (Modeled)

Computer models of the photometric consequences of errors in Y-base positions are being developed at the STSCI using the (ever-improving) FOS computer simulation. Dr. Ralph Bohlin will provide the results as an update to this section.

## 6.4 Pulse Coincidence Correction

- 6.4.1 Channel Deadtime The two 512 channel FOS sensors are multichannel photon counting detectors. Each of the separate 1024 channels has a 10.5 microsecond dead time circuit. That is, upon arrival of a channel event, that electronic channel is digitally gated off for 10.5 microseconds. This technique forces each channel to have the same input to output counting characteristic.
- 6.4.2 <u>Pulse Pair Calibration</u> The vital first element of the FOS calibration effort is the accurate removal of this digicon electronics nonlinear pulse counting characteristic. We assume the standard formula:

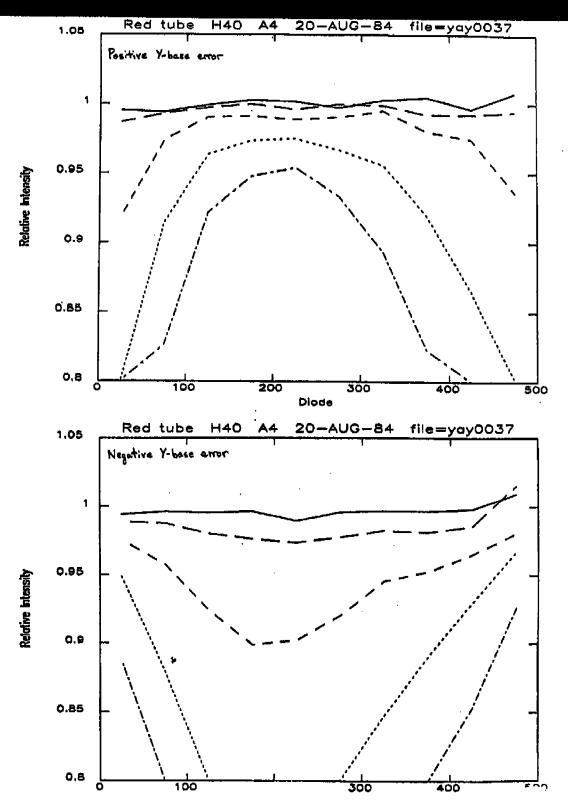


Figure 6.3.3-2. The effects of error in Y-position from continuum spectra on the red tube. The lines are coded at in Figure 6.3.3-1. The noise level is 1-2%, because the total counts per bin is higher than that of the emission line spectra.

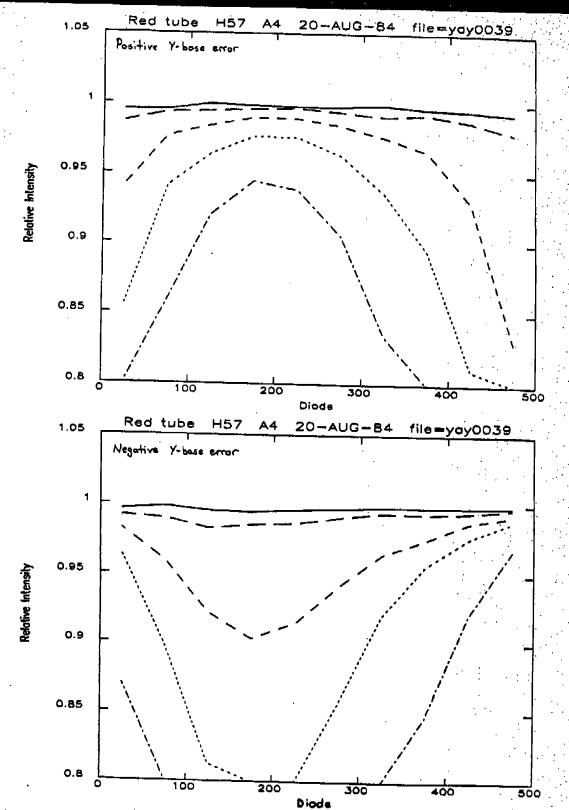


Figure 6.3.3-3. The effects of error in Y-position from continuum spectra on the red tube. The lines are coded as in Figure 6.3.3-1. The noise level is 1-2%, because the total counts per bin is higher than that of the emission line spectra.

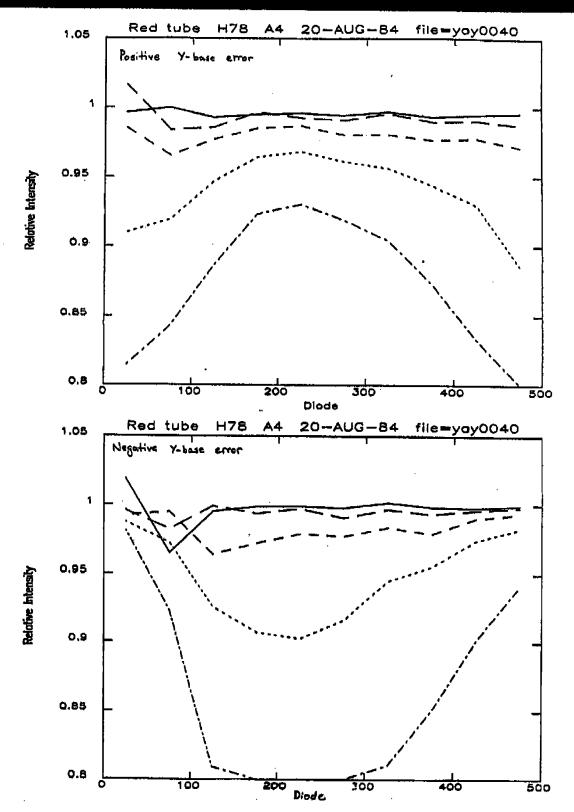


Figure 6.3.3-4. The effects of error in Y-position from continuum spectra on the red tube. The lines are coded as in Figure 6.3.3-1. The noise level is 1-2%, because the total counts per bin is higher than that of the emission line spectra. -233-

$$Y = \frac{X \exp(-t_1 X)}{1 + t_2 X \exp(-t_1 X)}$$
 (6.4.1)

where X --> true input count rate

Y --> apparent output count rate

t ~-> recovery time

t --> dead time

Pulse paired correction calibration involves the optimum determination of  $t_1$  and  $t_2$  and assessing the accuracy of the standard formula (6.4.1) to correctly linearize digicon data.

6.4.3 FOS Counting Data - During FOS calibration, the input light rate was varied over a range of three magnitudes by sequencing through the 11 separate apertures. Table 6.4-1 gives the relative ratio of the various apertures to the B1 aperture. The ambient STOS serves as the constant light input to the FOS apertures.

The FOS calibration data is plotted in Figure 6.4-1. The continuous curve in Figure 6.4-1 is a least squares fit of Equation (6.4.1) to the data, assuming  $t_2=10.5~\mu s$ . The optimum value of  $t_1$  is -0.17  $\mu s$ . Except for real counting rates  $X>10^5$  counts/sec, the fit is excellent, as displayed in Figure 6.4-2. The FOS maximum counting requirement is  $10^5$  counts/sec. Any deviations from the dominant  $t_2=10.5~\mu s$  dead time can be attributed to errors in the aperture sizing at this accuracy level.

6.4.4 <u>Procedure for Pulse Pair Calibration</u> - For input rates below  $10^5$  counts/sec. Equation (6.4.1) can be used to establish the true count rate.

In practice, we have implemented pulse pair calibration by a procedure involving interpolation of a look up table. This also

TABLE 6.4-1

Aperture	Relative Size (Aperture/B1)
A4*	0,111
A3	0.898
B2	1.00
B1	3.19
A2	4.07
C2	6.03
BS	13.4
C4	14.3
C1 :	17.0
09	42.8
A1	104.1

<sup>\*</sup>Results from this aperture have not been considered, since effective size is difficult to measure due to diffraction effects.

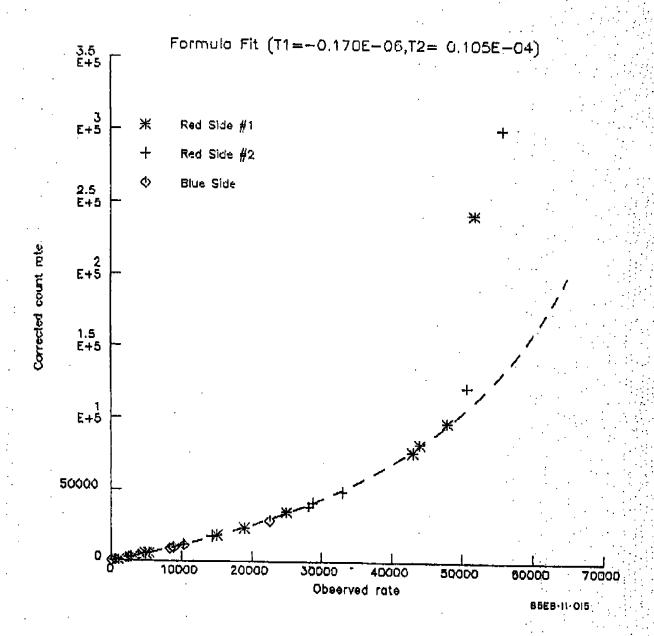


FIGURE 6.4-1

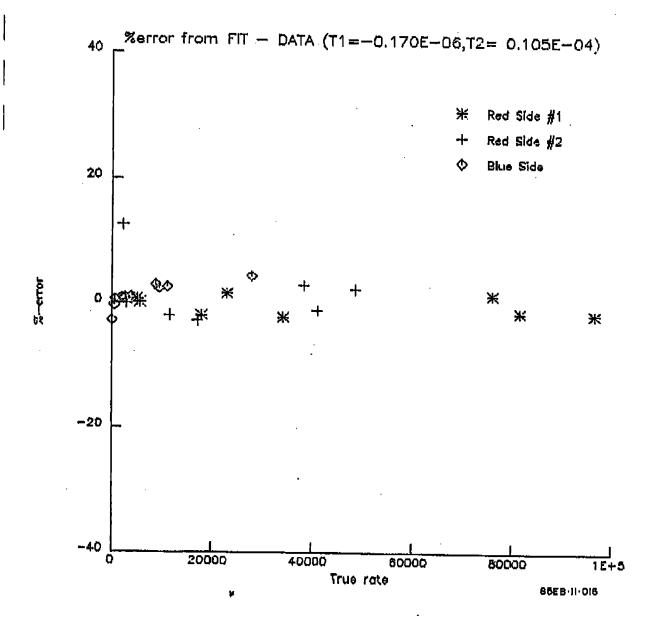


FIGURE 6.4-2

allows for pulse pair calibration at input rates  $X > 10^8$  counts/sec, where Equation (6.4.1) does not apply.

### 6.5 Background

### 6.5.1 Dark Count Rate

A series of tests have been conducted with the FOS flight and spare detectors, during various stages of fabrication and assembly, and under a variety of environmental operating conditions.

At the simplest level, dark count data were obtained for samples of various diodes, near the desired operating voltage, for all digicon phototubes built at EVSD. These data represented a portion of the overall tube data used for the selection of the best tubes for flight. The data were obtained at an operating temperature of approximately -10°C, where the dark count spec (0.002 counts/sec/diode) must be verified. A sample of such data are included here.

At the next level, the digicon tubes are encapsulated, wired, and fully integrated into a single operating unit which incorporates the structural, thermal, and HV interfaces along with the permanent magnets, the deflection coils, and the charge amplifier hybrid boards.

Dark count data were obtained with these fully assembled flight detectors to verify the dark count specification compliance at this integration stage. A description of the test set up, operating conditions, and sample results are also included here.

Finally at the systems level, where the detector assemblies are integrated with all other operating components of the FOS, dark count data were obtained to demonstrate systems specification compliance. These data were obtained under a wide variety of operating conditions and environments, primarily during ambient calibration and Thermal Vacuum Cycling and Thermal Vacuum UV calibration of the FOS. Again, samples of these test data are described and included here to illustrate systems level compliance with the specifications.

Overall, the test program at all levels demonstrates that the FOS does indeed perform at very low noise levels, well within its performance specifications.

### 6.5.1.1 Digicon Dark Count Data

Test data were obtained by Electron Vision Systems Division (EVSD) for the digicon vacuum phototubes upon completion of their fabrication. Sample test data were obtained for selected diodes, deflected to various photocathode locations. These data are shown in Figure 6.5.1.1-1 and were obtained with an operating voltage of 22 kV at a temperature of -11°C. The integration time interval was 1000 sec. This data sheet is for flight tube F-7, and represents the basic information to screen out problem tubes. In the example shown, the average dark count is seen to be approximately 0.001 counts/sec/diode, within the acceptable specification level of 0.002 counts/sec/diode.

# 6.5.1.2 Detector Assembly Level Dark Count Measurements

These tests were performed in a Tenney Model T5-110250 Thermal Chamber. All measurements were made at ambient (room)

#### AT-3010

### BACKGROUND COUNT RATE

Data Sheet

Tube S/N: 1=05 # =-7

Test Conditions:

Tube Temperature -10

Photocathode Voltage \_\_\_\_

BACKGROUND COUNTS PER 1000 SEC PER DIODE

HV . off = ,0008 Average Over Entire Field . 00 123 @ - 11°C

(Specification Limit = 2 Max.) (.005/ @ +/5° C.

Y Position (mm)

								<u>-</u>					
1	Diode #	1	59	105	159	203	259	311	359	405	459	509	<b> </b>
ļ	<b>/+3.</b> 5	1	6	.0	_	0		0		2		3	
1*	0	ユ	ユ.	5	え	1	2	3	ス		0		
-* .	-3.5	0	O	.0	0	2	0	1.	1	0	0	2	
(	<b>(+3.</b> 5	2	8'	1	1	0	0	0	1	2	1	0	
2**	}	උ	2	0	1	0	0	0	0	2	ノ	0	
4 X	-3,5	1	4	ن	- 1	0	4	o	1	0	0	ス	
•	HV Off	1	5	o	0	0	2	0	0	0	0		
外步火	"Room" Temp.	5-	10	<u></u>	11	6	5	1	6	4	3	0	1

\*1 Diode #1 near edge of mask.

\*\*2 Diode #512 near edge of mask. \*\*\*3 Terip <u>+/5</u> °C.

Date: 8-3-8/

Engineering:

FIGURE 6.5.1.1-1

-240-

pressure. To maintain a dry environment and to avoid condensation on the detector, the chamber was purged with dry gaseous nitrogen throughout the operation, except for periods of high voltage operation. During high voltage operation, the detector was purged with Sulfur Hexafluoride to minimize the effects of high voltage breakdown in the tube's high stress front end.

The detector was mounted in a Permanent Magnet Focus Assembly (PMFA). An array of "pogo pins" was used to contact the detector's header bumps to the input of the SAI Digicon Test This facility contains twelve channels of charge Facility. sensitive amplifiers, twelve amplifier/discriminators, and a Hewlett-Packard 9845 computer for test control and data preprocessing. A test program in the computer software permits monitoring of all twelve channels simultaneously throughout the 1000 one-second integration periods used for each test. A green LED is included in the test facility to permit illumination of the cathode for aliveness checks. Aliveness checks were made before and after each set of data to verify proper detector and test facility operation.

It should be noted that the discrete charge amplifiers used in these tests, although of the FOS design, are somewhat noisier than the hybrid charge amplifiers used in the flight configuration to be flown on the FOS.

### 6.5.1.2.1 Digicon F-3 (S-20 Photocathode)

This was the first digicon to be tested. Room temperature baseline dark counts were made, recording four channels distributed across the diode array (some of the pogo pin contacts were

lost during thermal cycling). The average dark count was measured to be approximately 0.3 counts per second per pixel (see Table 6.5.1.2.1-1).

The chamber temperature was decreased to 0°C and the digicon cold soaked a minimum of four hours to permit full thermal stabilization. The average dark count was measured to be approximately 0.01 counts per second per pixel.

The chamber was returned to ambient temperature overnight and the room temperature dark count remeasured prior to beginning another cold cycle. This repeated the previously measured room temperature baseline of 0.3 counts per second per pixel.

The chamber temperature was slowly decreased to -10°C and the digicon cold soaked for a minimum of four hours. The average dark count was then measured to be approximately 0.003 counts per second per pixel.

The chamber was returned to room temperature. A final room temperature dark count was made before the digicon was removed from the chamber. This data averaged approximately 0.3 counts per second per pixel, repeating the room temperature baseline rate.

### 6.5.1.2.2 <u>Digicon F-5 (Bi-Alkali Photocathode)</u>

This, the second digicon to have its dark count measured, was by far the more difficult. Not only was the count rate from the UV photocathode lower, but we had more problems associated with noise in the test set-up during these measurements.

The average dark count at room temperature was found to be .0038 counts per second per pixel.

TABLE 6.5.1.2.1-1

### FOS DIGICON NO. F-3 DARK COUNT

(August 7, 1983)

Temp.	Diodo Va		<b>0</b>	L. //	_	_	
	<u>Diode No</u> .		coun	ts/1000	<u>sec</u> ,	AVG	<u>Dark Count</u>
20°C	105	287	306	376			
	256	347	342	306			•
	355	331	305	268			
	509	<u>360</u>	<u>391</u>	<u> 388</u>	·		
		1325	1344	1238			
	•	(.331)	(.3\$6)	(.309)		•	.926
O°C	105	12	9				
	255	11	11				
	355	10	4				
	509	8	12				
		41	36				
		(0.10)	(.009)				.0096
-10°C	105	. 1	1	2	4	4	
	255	3	4	1.	1	4	
	355	2	1	3	2	2	
	509	<u> </u>	6	4	<u>7</u>	7	
		12	12	14	14	17	
		(.003)	(.003)	(.0025)(	0042}	,0033	

At reduced temperatures, the dark count dropped to .0014 counts per second per pixel at  $+3^{\circ}$ C, and to an average of .00075 counts per second per pixel at  $-11^{\circ}$ C.

# 6.5.1.3 System Level Dark Count Measurements from T/V

Figure 6.5.1.3-1 summarizes all background counting rate data points taken during T/V II. The important feature to observe in this figure is that there is no significant increase in background rates for either detector with increasing voltage.

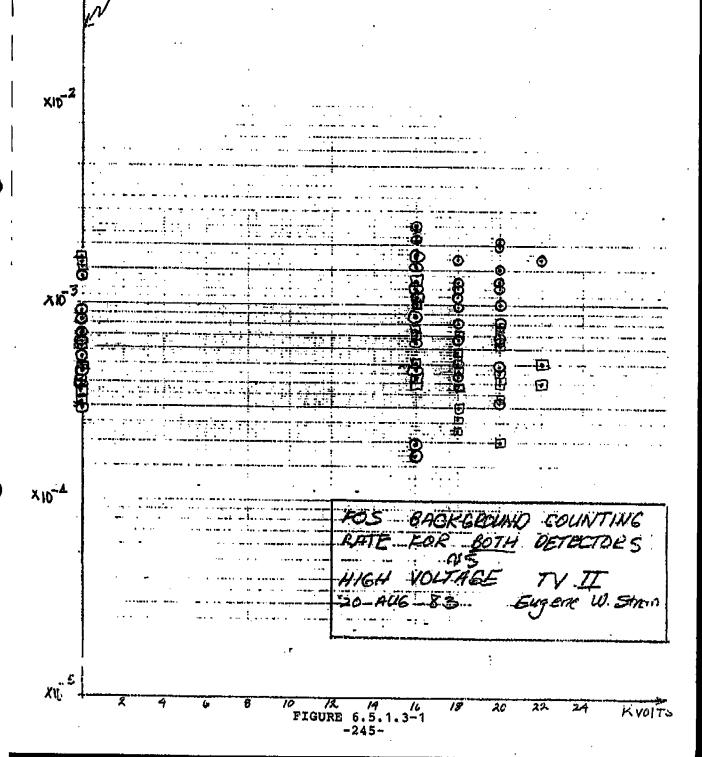
The background counting rate vs. temperature for the red and blue detectors are illustrated in Figures 6.5.1.3-2 and 6.5.1.3-3, respectively. There are no significant changes in the background rates with temperature. The corrected data result from eliminating counts from individual noisy channels (28 channels on the red side and 12 channels on the blue side) to get a background that is representative of the majority of pixels.

The background counting rates vs. voltage at constant temperature for the red and blue detectors are illustrated in Figures 6.5.1.3-4 and 6.5.1.3-5, respectively. Again, the important feature to note in these figures is that there is no significant increase in counting rates as the detector bias voltage is increased.

The background counting rate is less than the specification requirement of .002 counts/sec/diode for both detectors.

Figure 6.5.1.3-6 is a graph illustrating the background rate for each channel during the long background test for the blue detector. The baseline rate for this test is  $4 \times 10^{-4}$  counts/sec/diode. Channels 77, 146, 170, 177, 232, 237, 279, 422, 433,

# BACKGROUND RATE COUNTS/SECOND/DINCE



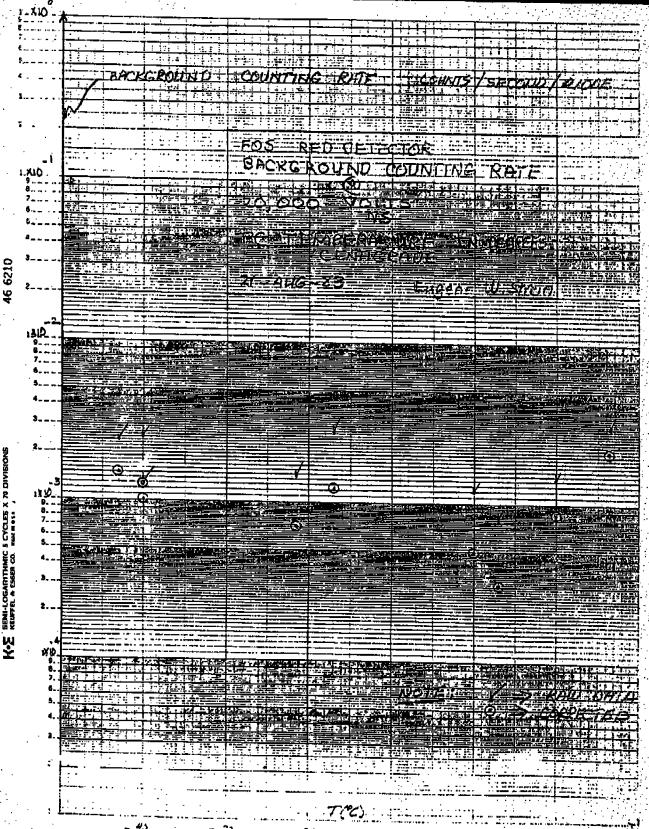


FIGURE 6.5.1.3-2

6.5.1.3-3 -24710

-248-

-249-

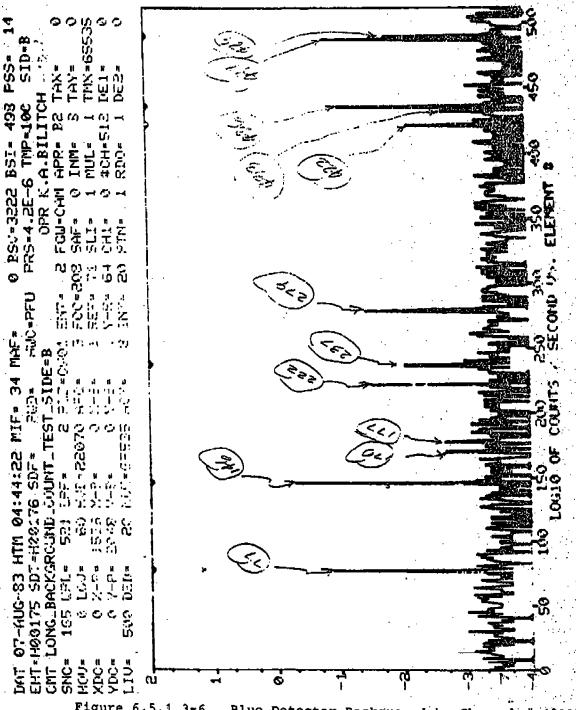


Figure 6.5.1.3-6. Blue Detector Background by Channel #, (022kV)

436, 487, and 489 exhibited background rates that were significantly above the baseline. The data were taken with the detector operating at 22 kV.

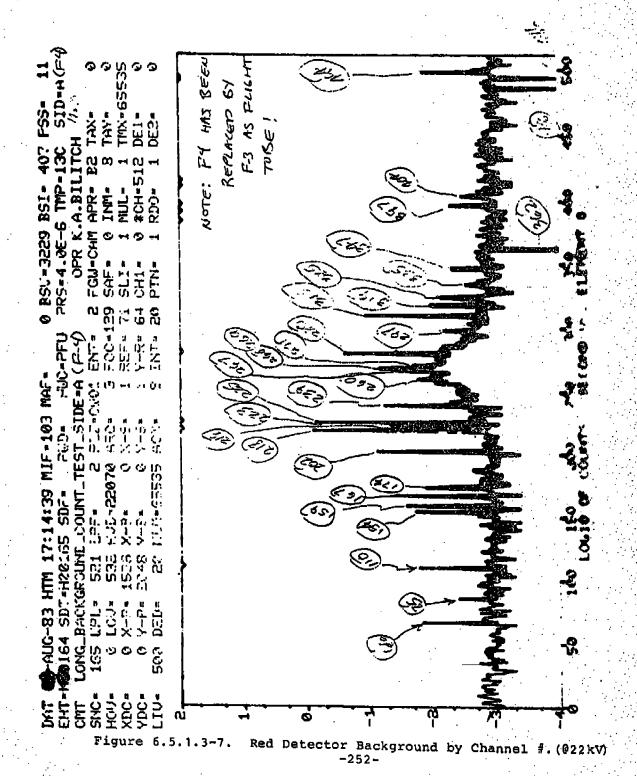
Figure 6.5.1.3-7 is a graph illustrating the background rate for each channel during the long background test for the red detector. The baseline rate for this test is 1.1 x 10<sup>-3</sup> counts/sec/diode. There are a large number of channels that have rates significantly above the baseline and there is a statistically significant feature at the center of the detector. The "hump" is believed to be additional evidence that the red detector is gassy. The enhancement in background rate at the center is believed to be caused by ions within the digicon. The data was taken with the detector operating at 22 kV. (NOTE: This F-4 red detector has since been replaced by F-3.)

Figure 6.5.1.3-8 is similar to Figure 6.5.1.3-7 for the Red Detector, except that these data were acquired at 20 kV. Note that the baseline rate is  $2 \times 10^{-3}$  counts/sec/diode for this test which is slightly higher than the data taken at 22 kV. The enhancement at the center of the detector is just noticeable here.

Figure 6.5.1.3-9 is similar to Figure 6.5.1.3-6 for the Blue Detector, except that the baseline rate is  $5 \times 10^{-4}$  counts/sec/diode for this test. Channels with counting rates significantly above the baseline are noted.

# 5.5.1.4 System Level Dark Count Measurements from T/V III

A number of the measurements performed during T/V II were repeated at T/V III. Some of these data are illustrated here, again to demonstrate overall system noise level specification



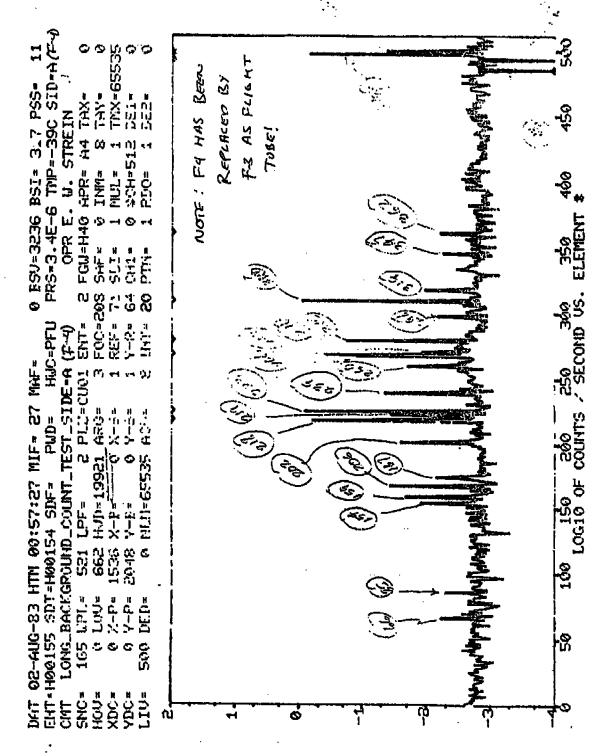


Figure 6.5.1.3-8. Red Detector Background by Channel #(@20kV) -253-

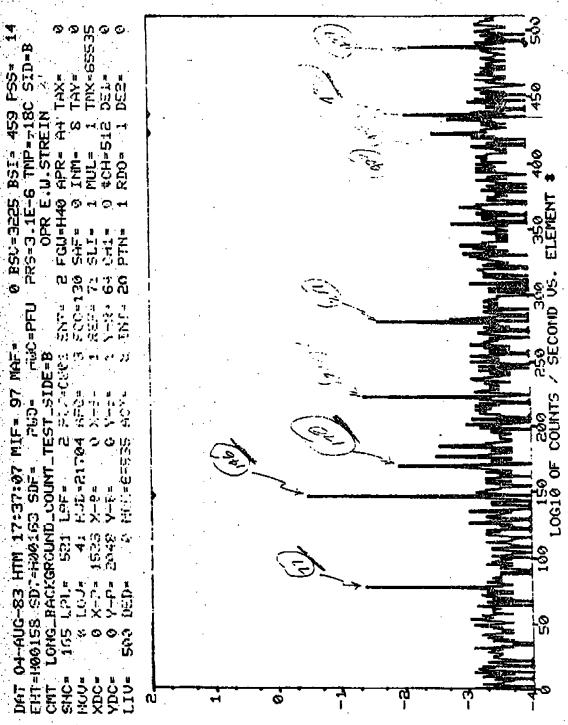


Figure 6.5.1.3-9. Blue Detector Background by Channel (@20kV)

compliance. Figures 6.5.1.4-1 through 6.5.1.4-3 illustrate dark count maps for the Blue Detector taken at various time intervals. The various noisy diodes which show the large noise spikes dominate the plots. These channels have been identified and will be treated to eliminate their contribution or replaced. The plots illustrate the repeatability and stability of the system level dark count within the specifications, after subtraction of the bad channel contributions.

Figures 6.5.1.4-4 through 6.5.1.4-7 illustrate a similar sequence of dark count plots for the Red Detector. Again, after subtraction of several bad channel contributions, the system dark count is stable and repeatable, and within specification.

Table 6.5.1.4-1 summarizes 100 sec integration interval dark count data as a function of temperature for the blue detector. Figure 6.5.1.4-8 illustrates a plot of some similar, but more limited, data obtained for the Red Detector.

### 6.5.2 Scattered Light

Scattered light levels are obtained by observing emission lines. At a distance of 50 diodes from the spectral lines, the scattered light ranges from 1 x  $10^{-5}$  to 4 x  $10^{-5}$  that of the peak count rates. Upper limits to scattered light at 2 diodes from the lines range from 7 x  $10^{-3}$  to 1 x  $10^{-2}$  the peak count rates. An anomalous feature on the blue tube with a periodic structure of 4 diodes is not understood.

A mercury discharge tube is observed and its spectra obtained with both tubes. Three narrow bandpass filters (~ 100  $\mathring{A}$  FWHM) are used to isolate bright lines at 2537, 3458, and 5461  $\mathring{A}$ .

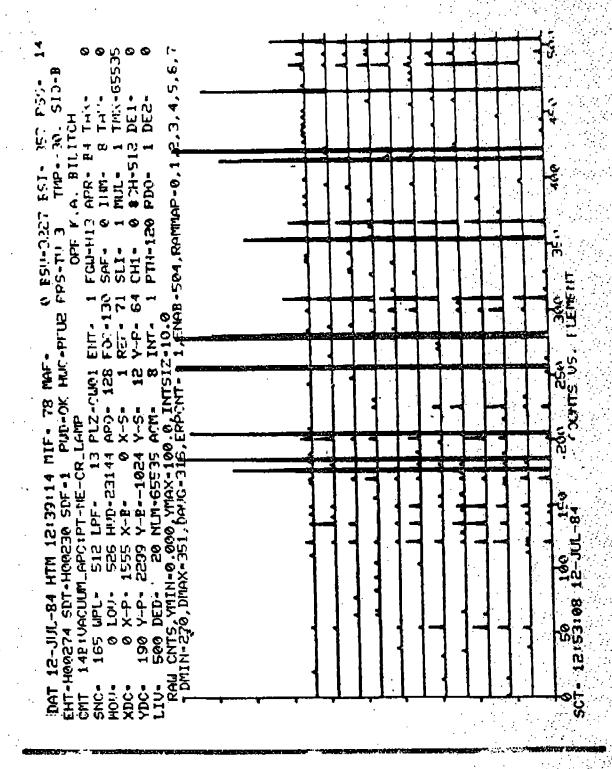


Figure 6.5.1.4-1. Blue Detector Dark Count Map from T/V III, t1 -256-

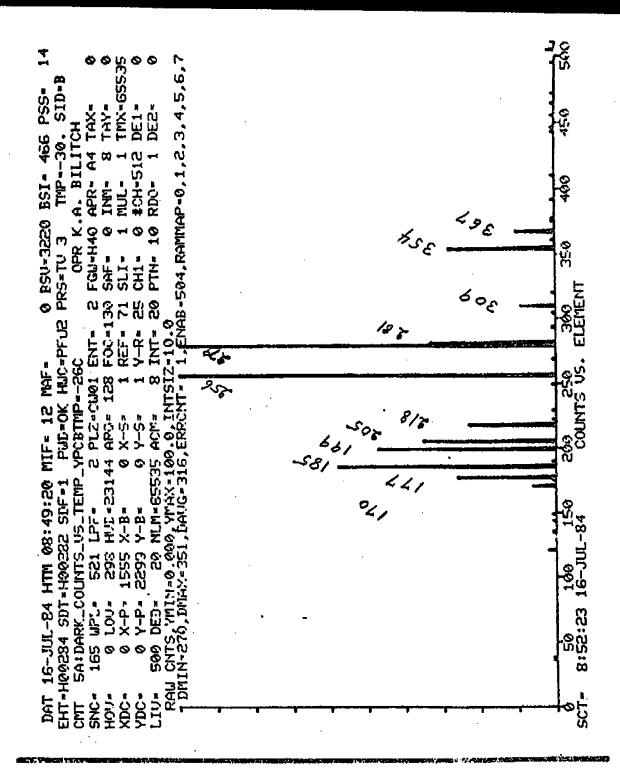


Figure 6.5.1.4-2. Blue Detector Dark Count Map from T/V III, t2 -257-

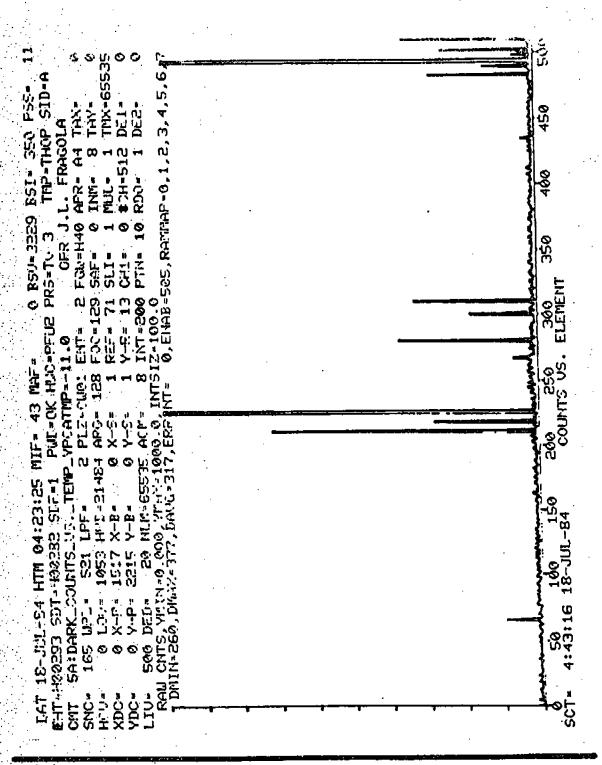


Figure 6.5.1.4-3. Blue Detector Dark Count Map from T/V III, t: -258-

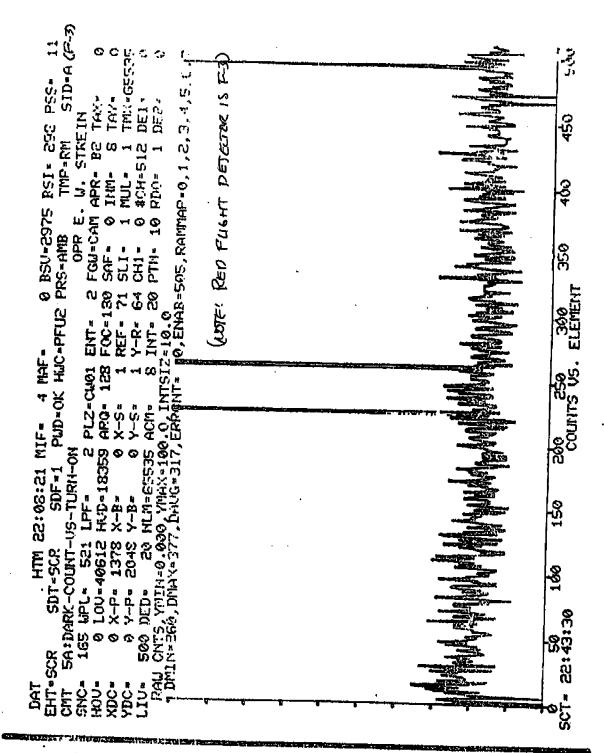


Figure 6.5.1.4-4. Red Detector Dark Count Map from T/V III, tl -259-

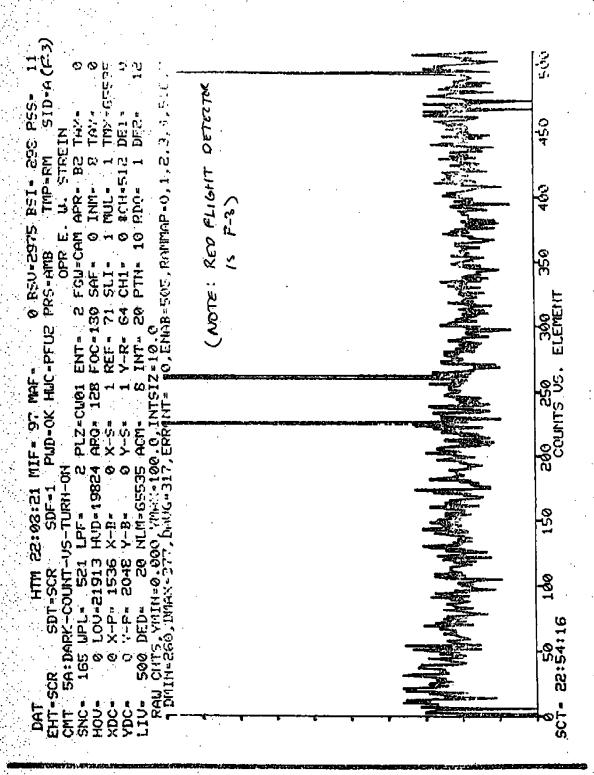


Figure 6.5.1.4-5. Red Detector Dark Count Map from T/V III, t2

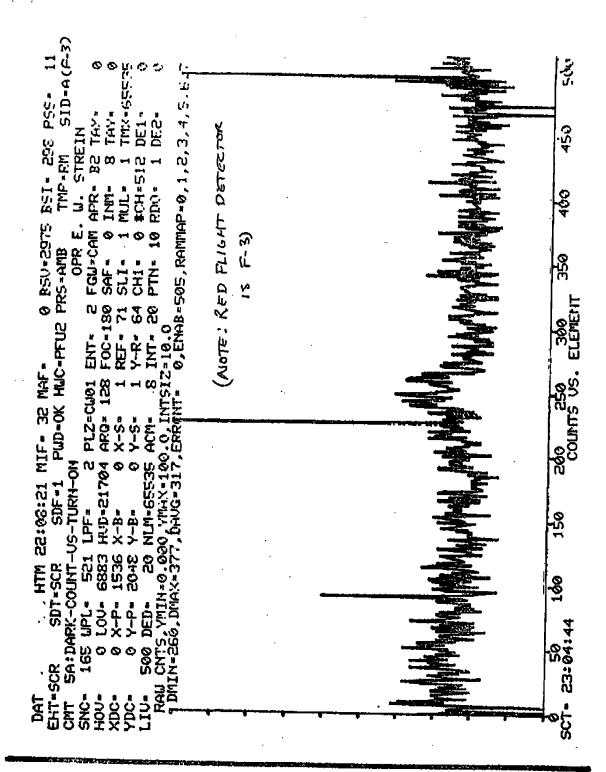


Figure 6.5.1.4-6. Red Detector Dark Count Map from T/V III, t4 -261-

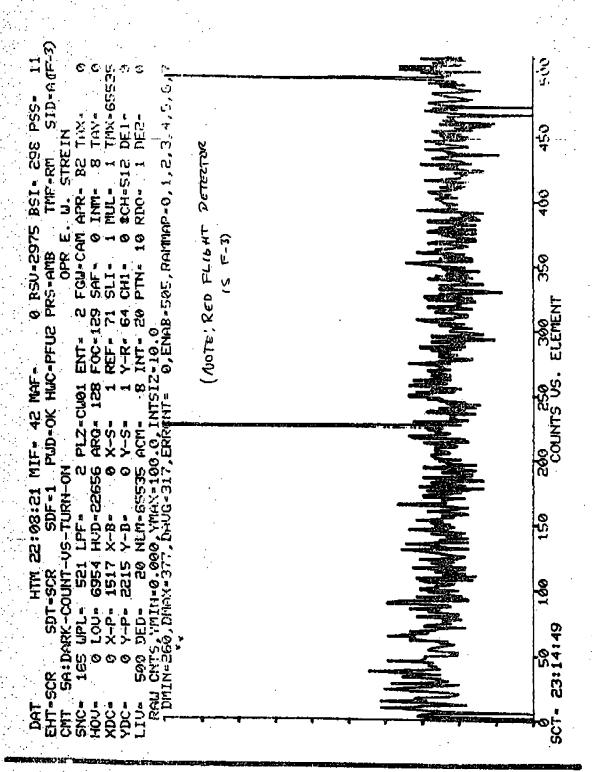


Figure 6.5.1.4-7. Red Detector Dark Count Map from T/V III, t4 -262-

TABLE 6.5.1.4-1

BACKGROUND RATE VS. TEMPERATURE, MONDAY, JULY 16, 1984

side = blue; 100 sec. integr. times

		,
TIME	<u> XPCBTMP</u>	BACKGROUND CT.RATE
07:52:34	-28C	8.06 (-4)
08:49:20	-26C	5.73 (-4)
10:58:37	-23.3C	2.48 (-4)
12:02:43	-22.7C	3.74 (-4)
13:13:24	~22C	4.63 (-4)
14:11:42	-21.3C	4.38 (-4)
15:29:14	-21.3C	5.71 (-4)
16:31:24	-20.7C	1.01 (-4)
17:24:03	-20.7C	6.85 (-4)
18:29:45	-20.0C	2.45 (-4)
19:29:38	-20.0C	5.25 (-4)
20:36:27	-20.0C	2.4 (-4)
21:20:47	-20.00	9.93 (-4)
22:17:57	-20.0C	4.61 (-4)
23:18:58	-20.0C	2.03 (-4)

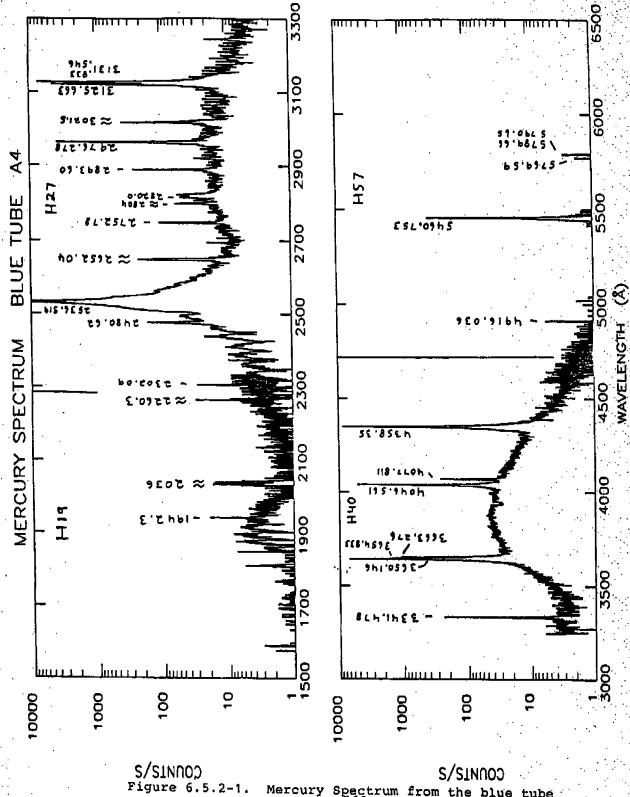
Poisson fit to low-noise diodes; in addition, there are non-thermionic noise spikes in typically 10 channels

Long and short exposures are obtained of each line using several gratings. Transmission ourves for the filters are determined using continuum lampa.

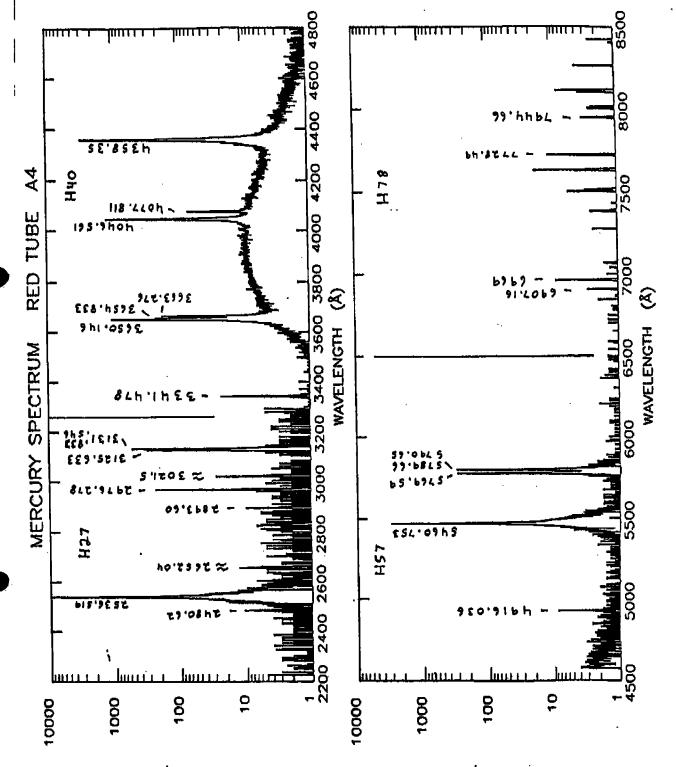
For each Digicon detector, four mercury spectra were obtained through the A4 aperture. Exposures of 5 seconds each were made with gratings H19, H27, H40, and H57 on the blue tube. On the red tube, exposures of 1, 15, 40, and 10 seconds were made with gratings H27, H40, H57, and H78, respectively. The concatenated spectra from the blue and red tubes are shown as Figures 6.5.2-1 and 6.5.2-2, respectively. Identified lines are labeled with their air wavelengths from the MIT wavelength tables. The blue spectra show identifiable lines from 1942 to 5791 Å and the red spectra show lines from 2481 to 7945 Å. The spectra from H40 for both tubes show a continuum from 3500 to 4600 Å. The bright UV line at 2537 Å on the blue tube has overflowed the 16-bit counter, but a shorter exposure through the 2537 A filter implies a peak of about 480,000 counts/sec.

Transmission curves for the three interference filters are obtained by taking the spectra of continuum sources with and without the filters (for example, see Figure 6.5.2-3). The expected spectra are obtained from the product of the mercury spectra and the transmission curves. Subtracting the expected spectra from the observed long-exposure filter spectra yields the scattered light (see Figure 6.5.2-4).

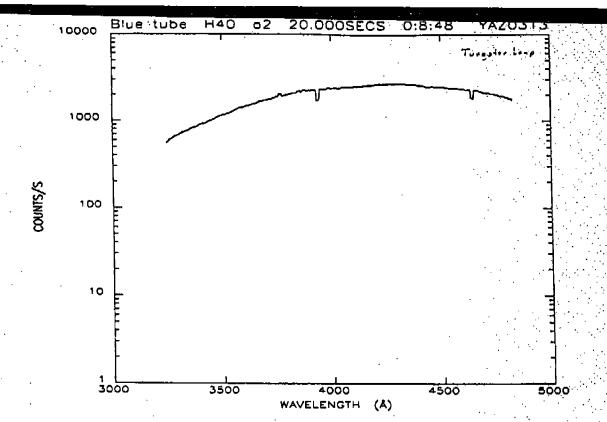
The short exposure spectra are used to estimate scattered light in the region 2 to 20 diodes from the mercury lines. Sub-traction of the expected spectra was not possible for these expo-

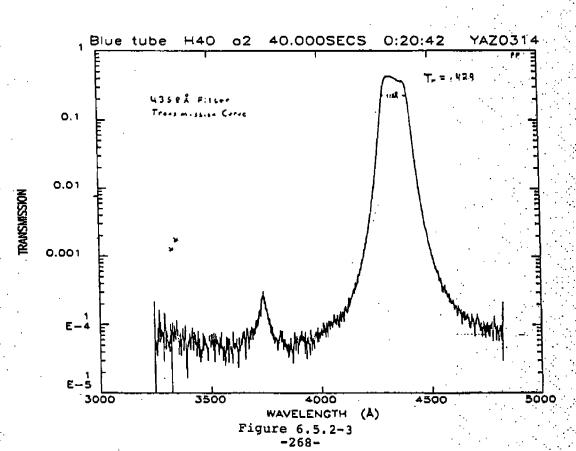


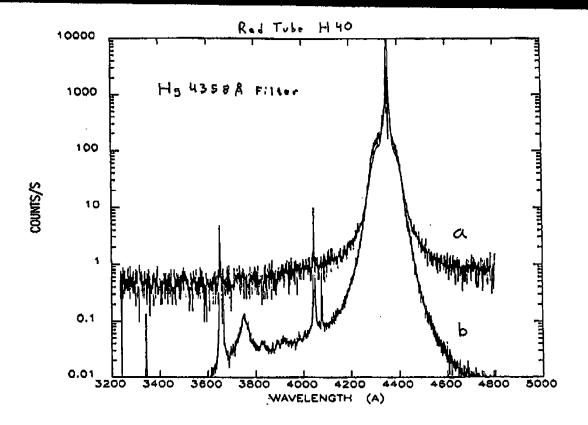
S/SINOOO Mercury Spectrum from the blue -266-

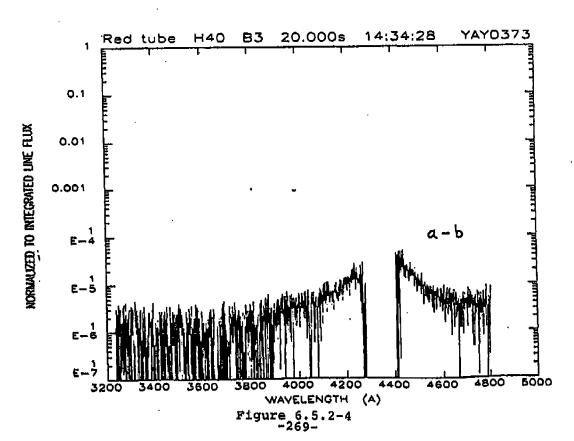


S/SINNOO S/SINNOO S/SINNOO Figure 6.5.2-2. Red Tube Mercury Spectrum as for Figure 6.5.2-1



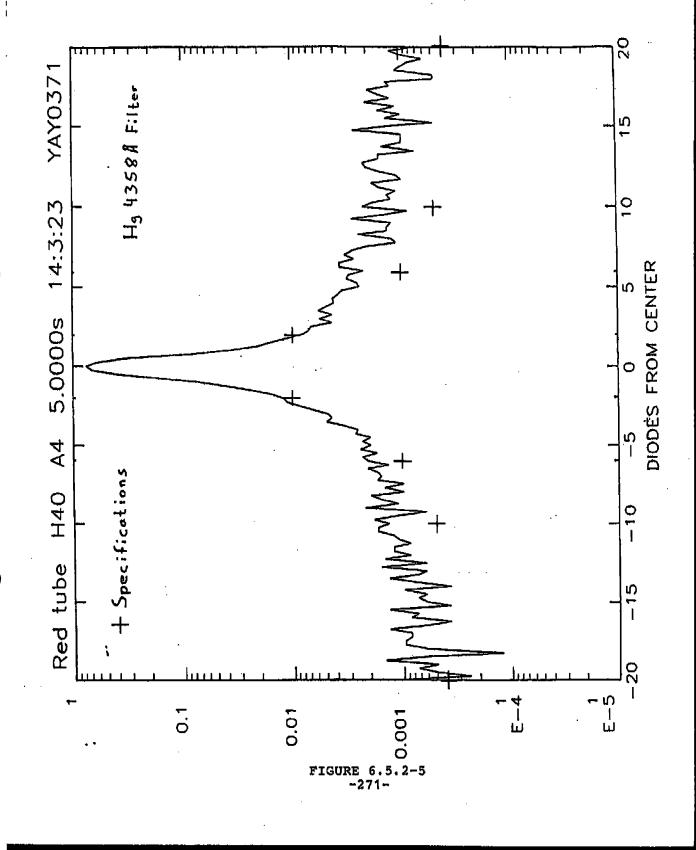


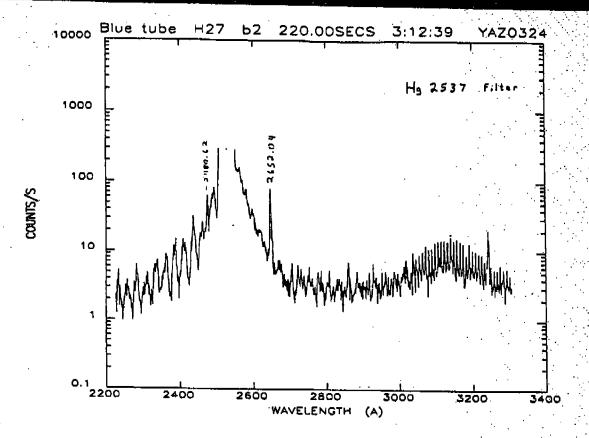


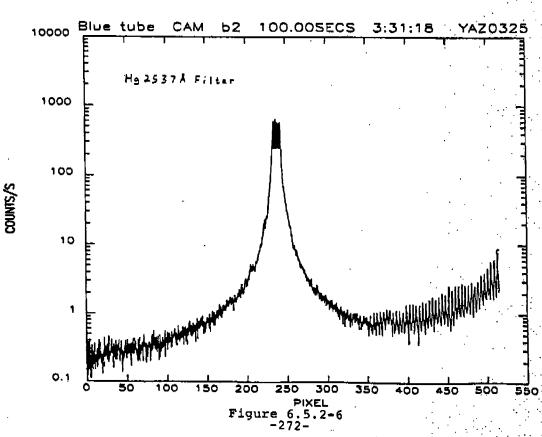


sures due to insufficient counting statistics. However, the spectra are upper limits on the scattered light. Figure 6.5.2-5 shows the profile of mercury line 4358 A observed on the red tube through the 4358 filter. At two diodes from the line center, the profile agrees with the requirements specified in the FOS design. From 6 to 20 diodes, the profile is greater than the specifications by about a factor of 2. Since the FWHM transmission of the filter is 38 diodes for H4O, the continuum in Figure 6.5.2-2, which is  $10^{-3}$  times the level of line 4358 Å, may dominate the profile, so that the actual scattering in the near wings may be well below the upper limits in Figure 6.5.2-5. The scattering in the far wings shown in Figure 6.5.2-4, bottom plot, is typical for both tubes and for gratings H27, H40, and H57. In all cases, the scattered light is well below the design specifications of  $2 \times 10^{-4}$  at 50 diodes and 1 x  $10^{-4}$  beyond 100 diodes from the line center.

In eight spectra from the blue tube, periodic features are present, examples of which are shown in Figure 6.5.2-6. In all eight cases, the period of the features is between 4.0 and 4.12 diodes and the distance from the emission line approximately 270 diodes. Since the features appear on spectra taken with H27, H40, L15, the prism, and the camera mirror, with filters 2537 and 4358, the anomalies seem to be independent of grating, filter, and wavelength, thus suggesting an origin within the Digicon. For the H27 spectrum in Figure 6.5.2-6, the maximum level of the anomaly is 2 x 10<sup>-5</sup> of the total count rate. No anomalous features are observed from red tube spectra. However, the thermionic







dark count of 0.33 counts/sec at ambient temperature may mask any such features on the red tube. Power spectrum analysis will be performed on red spectra to see if such anomalies do indeed exist.

## 6.5.3 FOS-Scattered Red Light (Red Tube)

The FOS is inherently susceptible to unwanted light, since it is a single pass spectrograph employing blazed, concave gratings, and possesses a wide spectral bandpass of sensitivity.

The laboratory data to be discussed here were obtained by placing several different (red-pass) absorbing filters in front of the (Tungsten) continuum source. The resulting absorption region was observed to measure the stray light contribution. This method is particularly relevant as it mimicks the observation of late type stars, which have rather steep cut-offs in their energy distribution in the ultraviolet.

The filters used were spare flight (Schott) units, designated OG830, GG395\*, WG295, and WG230. Their approximate transmission curves are shown as Figure 6.5.3-1.

The actual observed countrates (through aperture C3) are shown as Figure 6.5.3-2. For each filter, four different gratings (H19, H27, H40, and H57) were observed. The necessary paired-pulse corrections for countrates in excess of approximate-ly 2000 counts/sec are indicated with broken lines.

Of particular relevance are the observed counts using the H19 grating (1600-2350 Å). The four curves shown in that wavelength region correspond to each of the four filters used and the countrates increase from roughly 2 counts/sec for 06530 to about 8 counts/sec for WG230. This is easily understood as the conse-

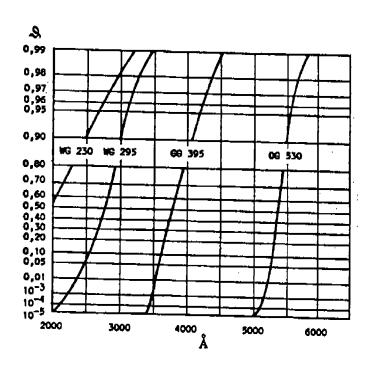


Figure 6.5.3-1

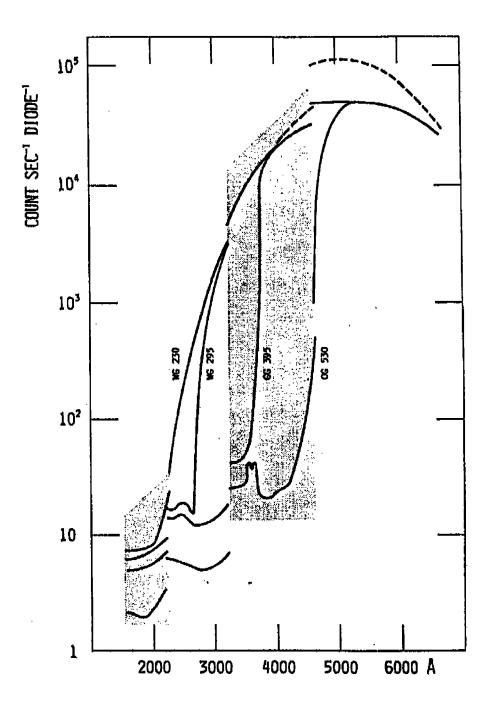


Figure 6.5.3-2

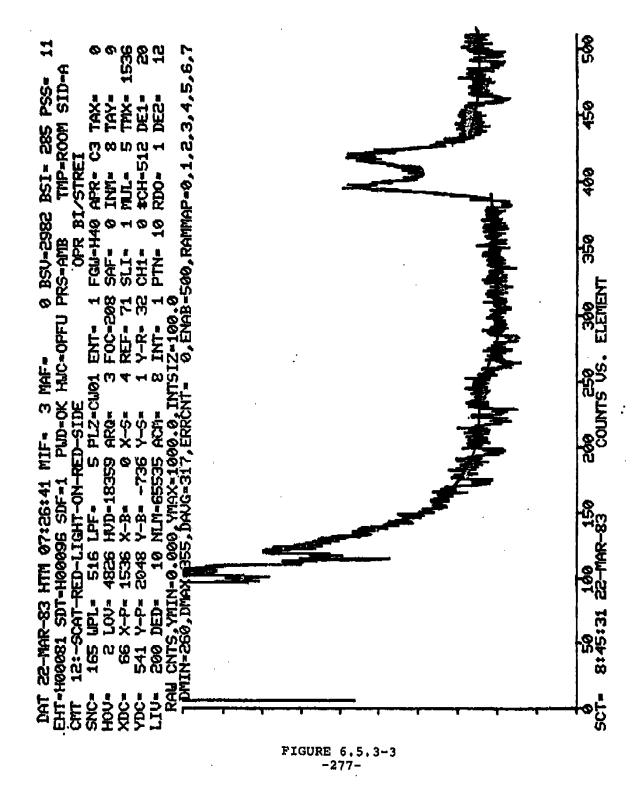
quence of the increase in bandwidth of the incoming flux from the Tungsten lamp. Note that although the Figure 6.5.3-2 is labelled in Angstroms, the measured counts are obviously from ("scattered") optical and near ultraviolet photons.

The effective bandwidth of the incoming light is limited at the long wavelength side by the cut-off of the photocathode of the (Red) tube which occurs around 7000 Å. On the blue side, the observed values for the cut-off are about  $4800^{++}$ , 3750, 2850 and 2300 Å, for 06530, GG395, WG295, and WG230, respectively. Note that although the increase of bandwidth in going from 0G530 to GG375 is rather modest in an absolute sense, the observed amount of scattered light (in H19) more than doubles. This is probably due to the increase of scattering efficiency for decreasing wavelength (e.g., Rayleigh scattering  $\alpha$   $\lambda^{-4}$ ) and suggests that most of the photons counted in H19 through the 0G530 filter have probable wavelengths of around 5500 Å, rather than the larger values present in the incoming flux.

The observed curves of Figure 6.5.3-2 are thus generally well understood, at least in a qualitative sense. Exception should be made for the unidentified feature around 3550 Å (H40-grating). The width of this feature is about 150 Å and exhibits a central "inversion". The observational data in this wavelength region are attached as Figure 6.5.3-3.

<sup>\*</sup>There is some confusion with respect to this filter. The FOS Scientists Notebook calls this filter GG395, whereas Harms et al. in NASA CP-2244 quote GG375.

<sup>\*\*</sup>This observed cut-off is somewhat surprising as the Schott-catalog (see Figure 6.5.3-1) claims that the transmission in this filter should be down to  $10^{-5}$  at 5000 Å.



## 7.0 STABILITY AND REPEATABILITY

## 7.1 High-Voltage, Thermal, and Aperture Mechanism Effects

Operational constraints allow the high voltage to be on for only one FOS detector at a time. After turning on the high voltage, the electric field in the Digicons requires time to achieve equilibrium. For the first 30 minutes after high voltage turn on, large image shifts (on the order of 1 micron per minute) are present. This motion decreases to less than 2 microns per hour after 60 minutes. Studies of the motion for a longer period using data taken to measure aperture repeatability, showed a small (1.2 micron/hour) drift up to five hours after turn on. However, this drift may be due to temperature effects.

The aperture repeatability data demonstrated that the aperture position on the photocathode was repeatable to within one micron.

During the transition from cold to hot operate, a motion of the spectrum along the diode array of 32 microns was measured on the blue side for a change of 5°C in the optical bench temperature.

To monitor the image stability after high voltage turn on, a series of platinum-neon spectra was taken using grating H27. Data were acquired for both detectors for a period of one hour after high voltage turn on. Data were taken at both cold and hot operate temperatures. To monitor the image stability for a longer period, these data were supplemented with data taken to measure the aperture wheel repeatability. These later data were

of 7 dispersers. With the additional motor step, the standard deviation was found to be 3.2 microns, indicating an average improvement of more than a factor of 2! The improvement ranged from "20% to a factor of nearly 3.5, depending on the disperser selected. The full range of minimum to maximum X-position over the 9 trials for each of the tested dispersers is plotted in Figure 7.2.2-1, along with some previous measurements made in the same manner. The mean spread has been reduced by approximately a factor of 2, from 21 microns for the October 1984 measurements to 10.2 microns.

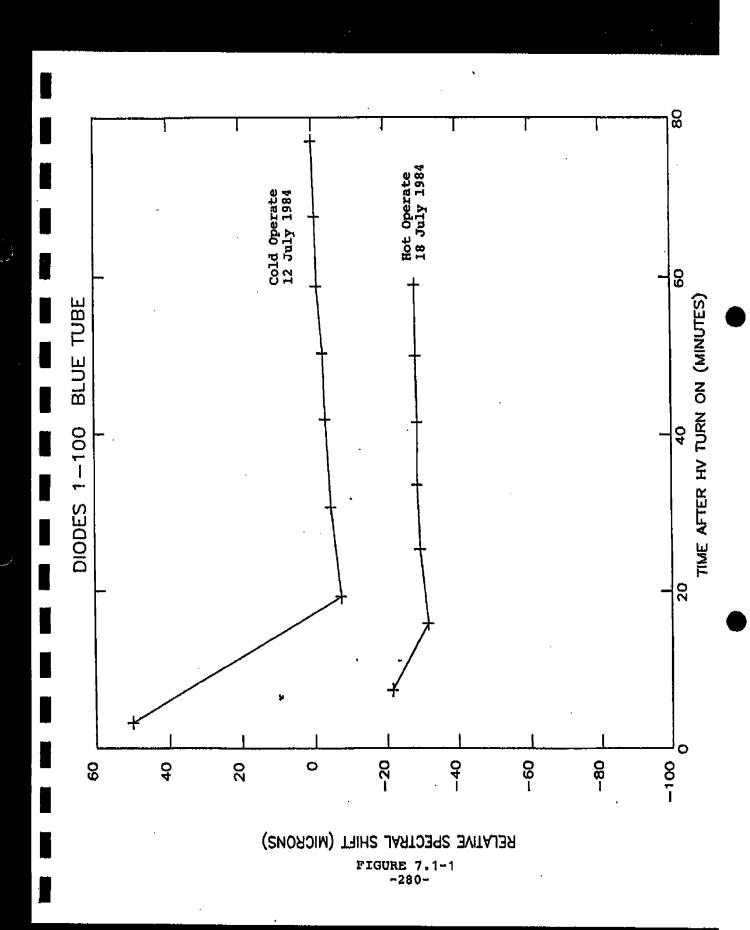
Although these data indicate that a noticeable improvement in the FGW repeatability is likely to result from the simple addition of a single step command following each wheel positioning, we note that even the performance level measured in the March 6 test is not entirely satisfactory, particularly for target acquisition into the smallest, 0.1 arcsecond (A4) apertures. The "10-micron spread in the (camera mirror) aperture image location, inferred from the average behavior of the 7 dispersers tested, remains a large fraction of the (14-micron) A4 aperture image size. For a substantial percentage of A4 acquisitions, the target will suffer significant decentering. drifts, which may now dominate the non-repeatability, will aggravate the situation. Furthermore, we have no assurance that the behavior demonstrated in our tests is applicable to the camera mirror, which was not actually measured, or that the improvement noted in the X direction will be reflected in the Y axis, since no correlation between X and Y errors was present in

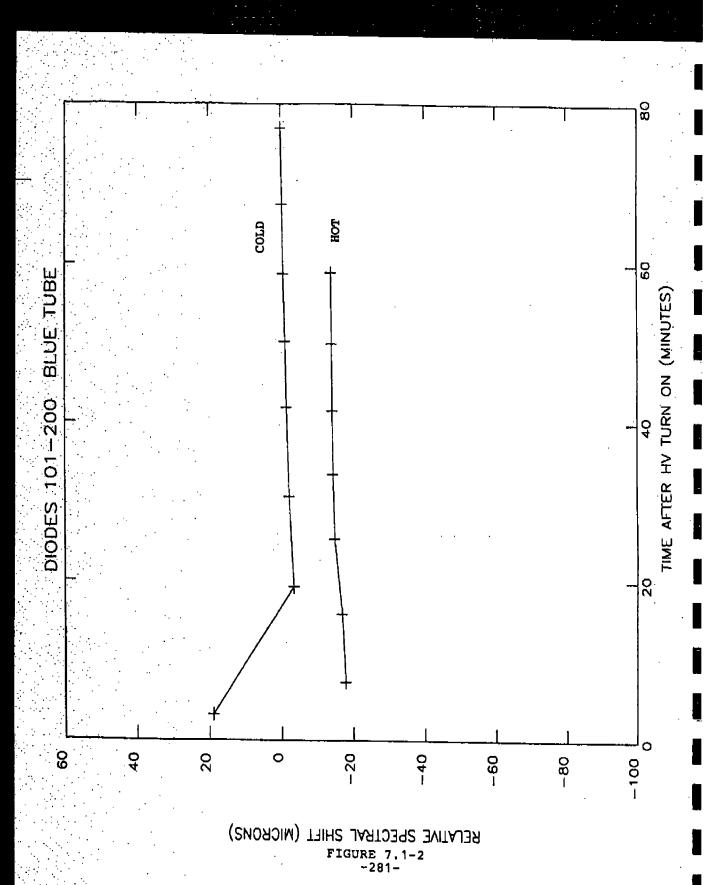
only available on the red side and extended to five hours after high voltage turn on.

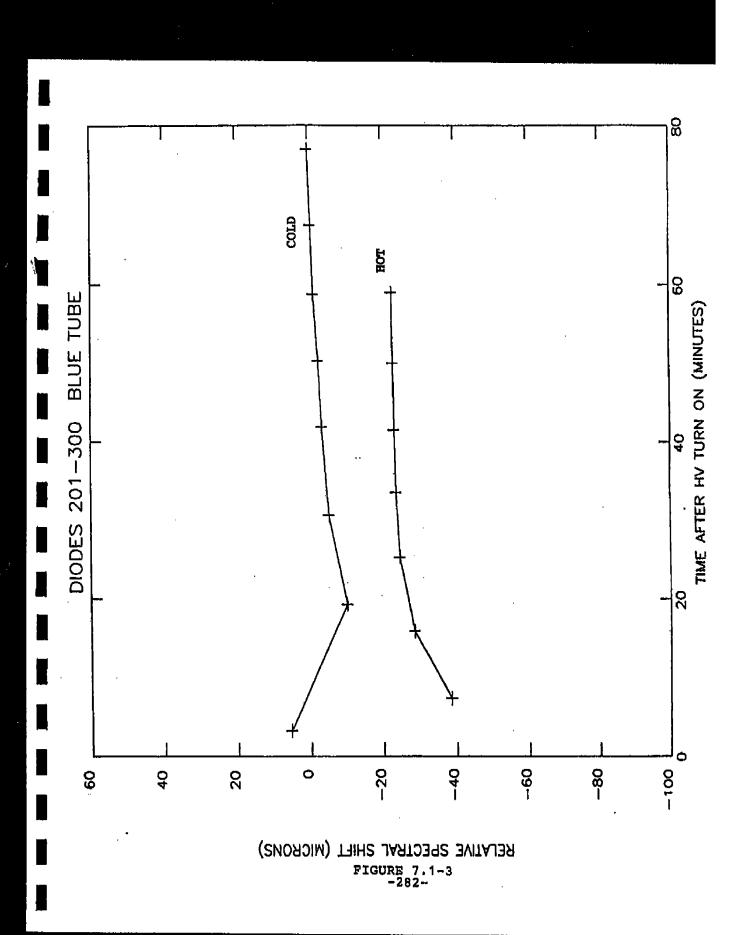
Motion between spectra was computed using cross-correlation techniques. In order to analyze motion versus diode position, the data were separated into groups containing diodes 1-100, 101-200, 201-300, 301-400, and 401-500. Cross-correlation was done on each group separately. To obtain motion estimates below the 50-micron resolution of the pixels and the 12.5-micron data sample spacings, the minimum of the correlation matrix was found and a quadratic fit was done using that point and its two neighbors. The position of the minimum of the quadratic fit was used to give the relative spectral motion.

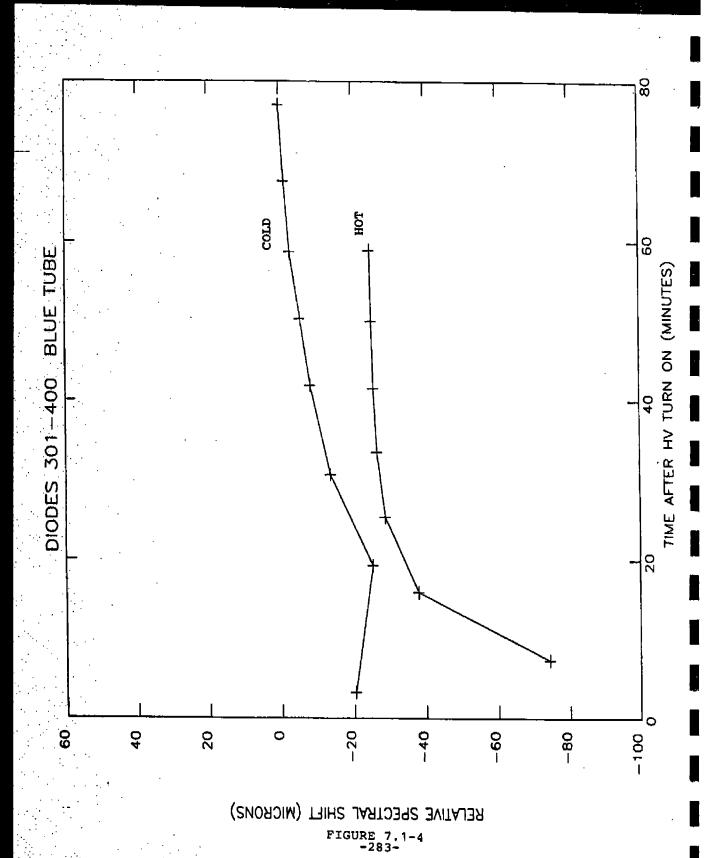
Figures 7.1-1 to 7.1-5 show the results for the blue tube in hot and cold operate. All shifts are relative to the last spectrum taken at cold operate. Large image motions are seen for the first 30 minutes with the largest shifts at the two ends of the diode array. The large offsets between the hot and cold operate spectra, seen 60 minutes after turn on, are due to the lack of perfect repeatability of the filter-grating wheel. Figures 7.1-6 to 7.1-10 show similar results for the red tube,

Figures 7.1-11 to 7.1-15 show an attempt to measure the high voltage settling after one hour using data taken to study the aperture wheel repeatability. All motion is relative to zero motion at 60 minutes after turn on. The aperture repeatability data were shifted in time by 45 minutes to match the previous data. After 100 minutes, all diodes moved in the same direction at a rate of 1.2 microne per hour. This small drift may be

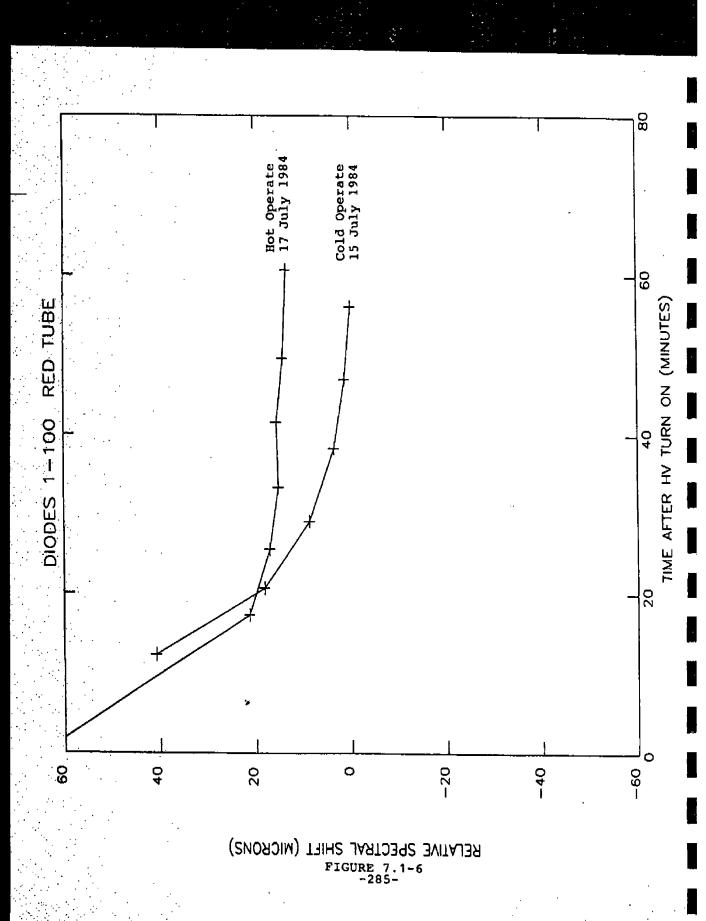


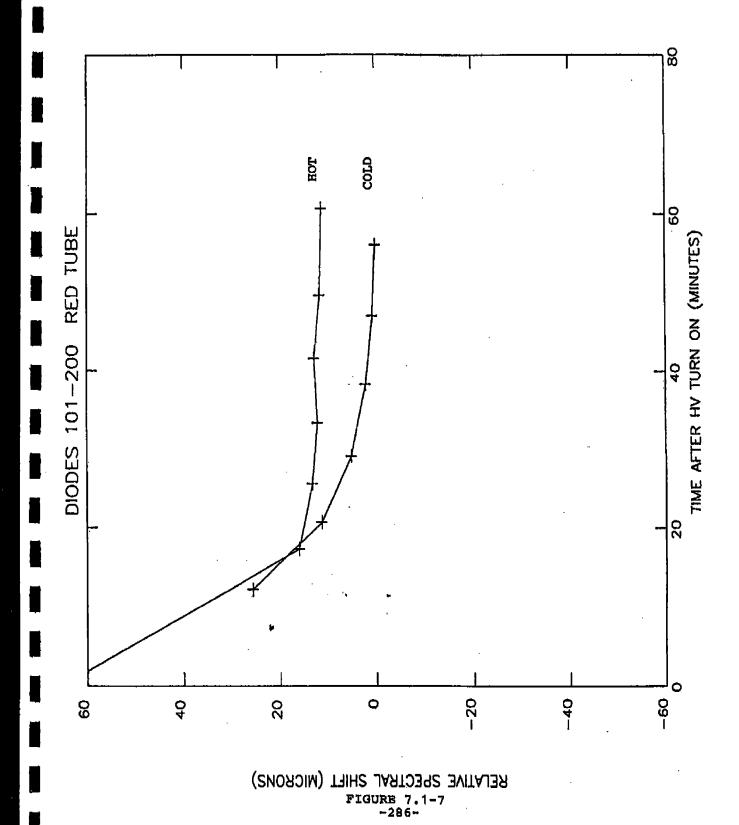


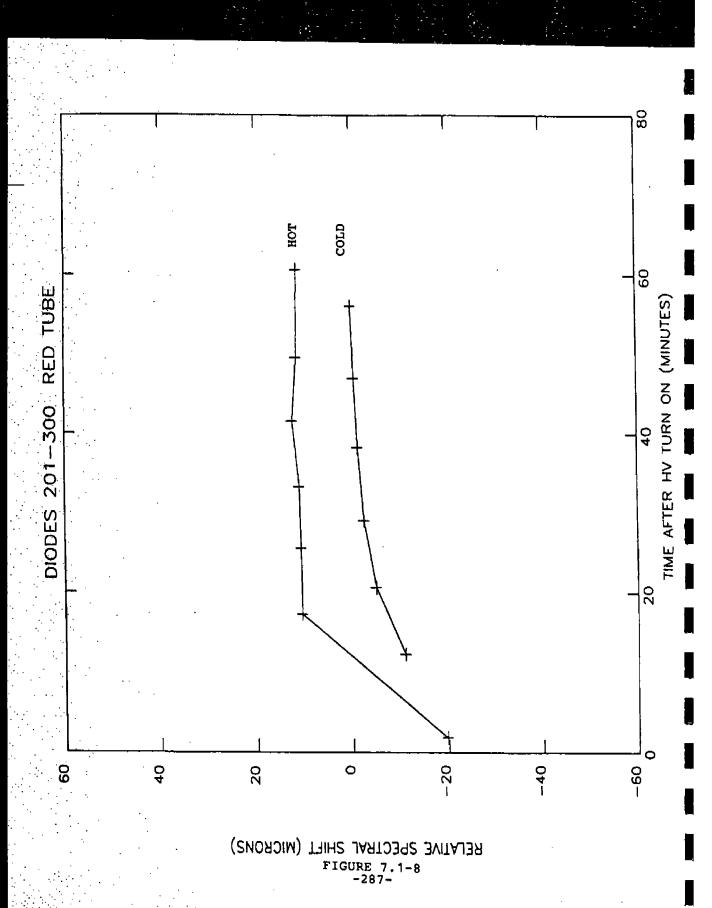


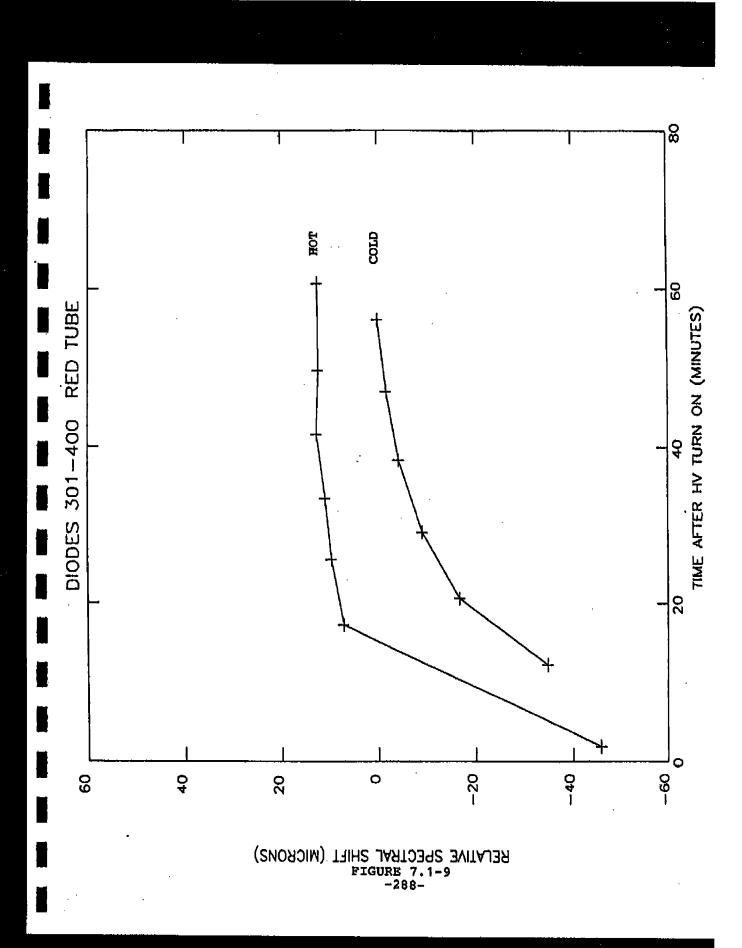


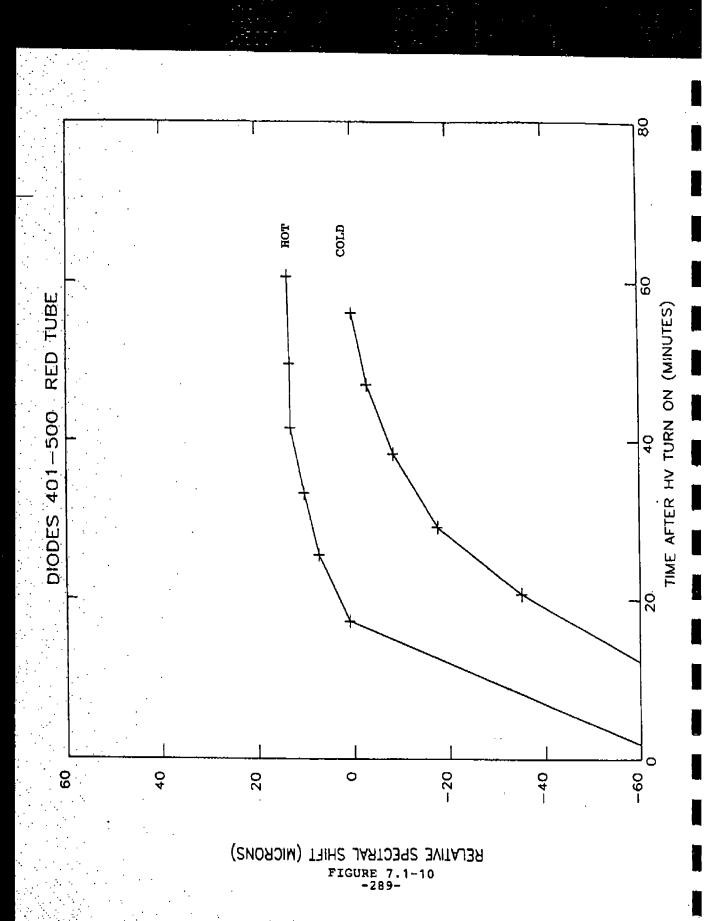
စ္ထ COLD HOT 8 BLUE TUBE TIME AFTER HV TURN ON (MINUTES) DIODES 401-500 4 9 5 20 -20 -40 -100 09--80 0 SELATIVE SPECTRAL SHIFT (MICRONS)
5-1.2 BARDIA
5-1.5 BARDIA

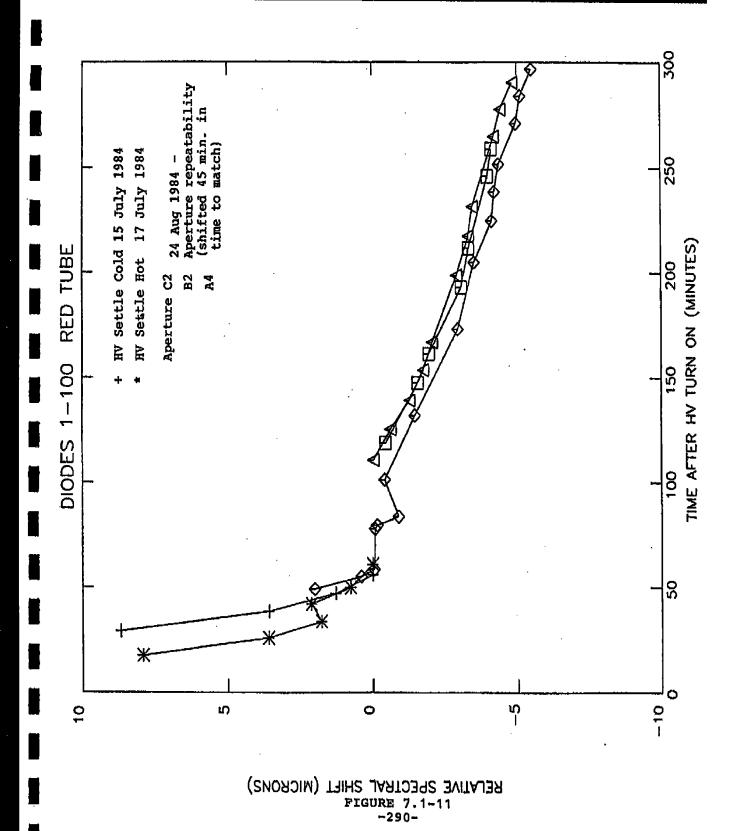












DIODES 101-200 RED TUBE TIME AFTER HV TURN ON (MINUTES) Ŋ KECATIVE SPECTRAL SHIFT (MICRONS)
-162-162-

